



# Manipulation of Lightwave Through Coordinate Transformation

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Ariadna bidder code: **21290**





# Project plan

We are here!



Tasks \ Time (mth)	0.5	1	1.5	2	2.5	3	3.5	4
Cylindrical cloak: Simplification								
Cylindrical cloak: Effect of order number								
<i>Electromagnetic Concentrator</i>								
Arbitrarily-shaped cloak								
Project report								



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Content	Time
0. Introduction	8
1. Cylindrical cloak: Simplification	15'
2. Cylindrical cloak: Effect of order number	6'
3. Electromagnetic Concentrator	6'
4. Arbitrarily-shaped cloak	6'
5. Conclusion	4'
Total:	45'



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# Introduction



Scattering → visibility

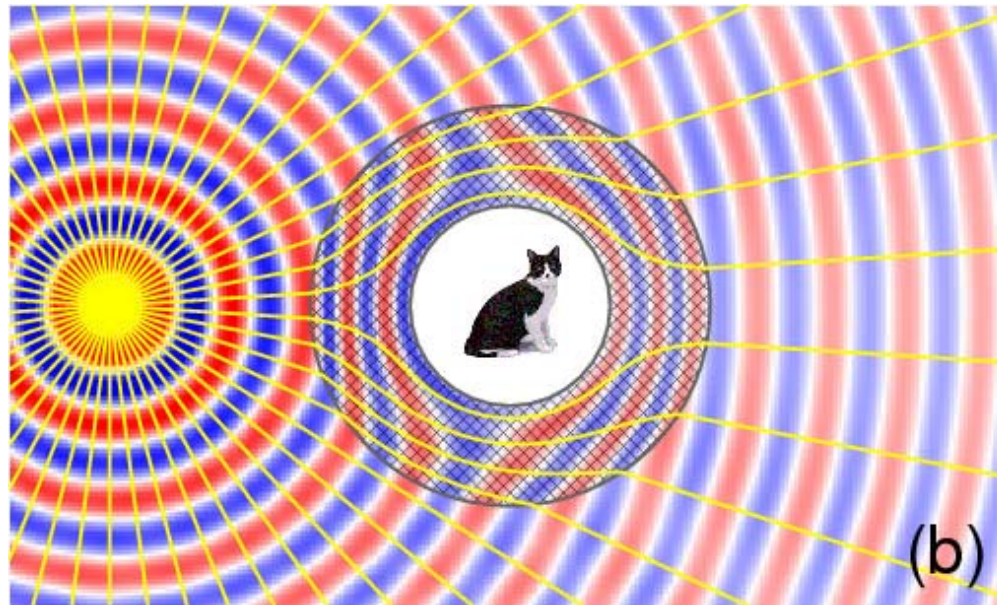


? → invisibility

Founder



# Coordinate transformation approach



1. J. B. Pendry, D. Schurig, and D. R. Smith, Science **312**, p.1780 (2006).
2. U. Leonhardt, Science **312**, p. 1777 (2006).



# Coordinate transformation

Original Cartesian coordinate  $(x, y, z)$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}, \quad \nabla \times \mathbf{H} = \frac{\partial \mathbf{D}}{\partial t} + \mathbf{j}, \quad \nabla \cdot \mathbf{D} = \rho, \quad \nabla \cdot \mathbf{B} = 0.$$

$$\mathbf{D} = \epsilon_0 \bar{\bar{\epsilon}} \cdot \mathbf{E}, \quad \mathbf{B} = \mu_0 \bar{\bar{\mu}} \cdot \mathbf{H},$$

Jacobian transformation matrix

New coordinate  $(q_1, q_2, q_3)$

$$x = f_1(q_1, q_2, q_3), \quad y = f_2(q_1, q_2, q_3), \quad z = f_3(q_1, q_2, q_3).$$

$$\Lambda = \begin{bmatrix} \frac{\partial x}{\partial q_1} & \frac{\partial x}{\partial q_2} & \frac{\partial x}{\partial q_3} \\ \frac{\partial y}{\partial q_1} & \frac{\partial y}{\partial q_2} & \frac{\partial y}{\partial q_3} \\ \frac{\partial z}{\partial q_1} & \frac{\partial z}{\partial q_2} & \frac{\partial z}{\partial q_3} \end{bmatrix}.$$

The Maxwell equations can take the invariant form as

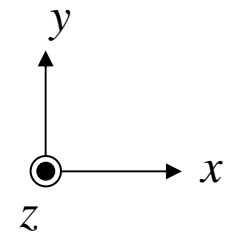
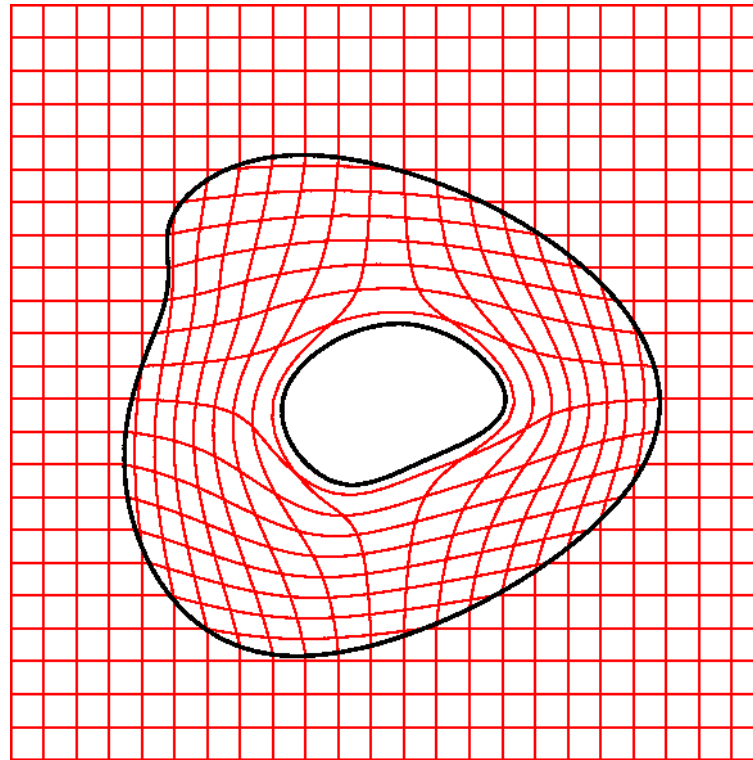
$$\nabla_q \times \hat{\mathbf{E}} = -\frac{\partial \hat{\mathbf{B}}}{\partial t}, \quad \nabla_q \times \hat{\mathbf{H}} = \frac{\partial \hat{\mathbf{D}}}{\partial t} + \hat{\mathbf{j}}, \quad \nabla_q \cdot \hat{\mathbf{D}} = \hat{\rho}, \quad \nabla_q \cdot \hat{\mathbf{B}} = 0$$

with

$$\hat{\bar{\bar{\epsilon}}} = \det(\Lambda)(\Lambda)^{-1} \bar{\bar{\epsilon}} \Lambda^{-T}, \quad \hat{\bar{\bar{\mu}}} = \det(\Lambda)(\Lambda)^{-1} \bar{\bar{\mu}} \Lambda^{-T},$$

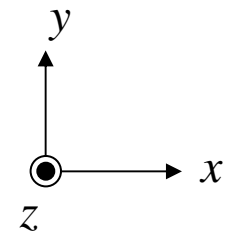
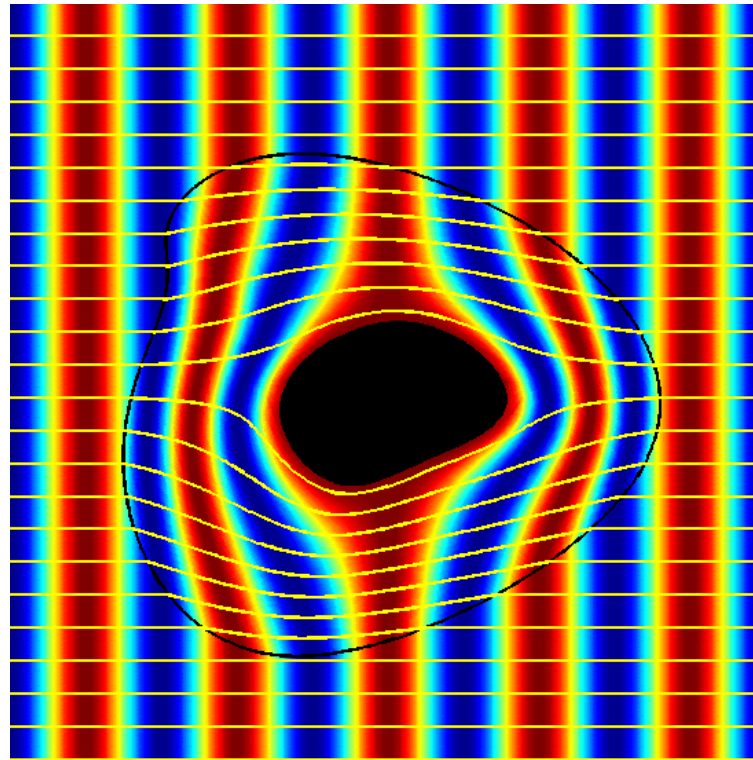


# Coordinate transformation





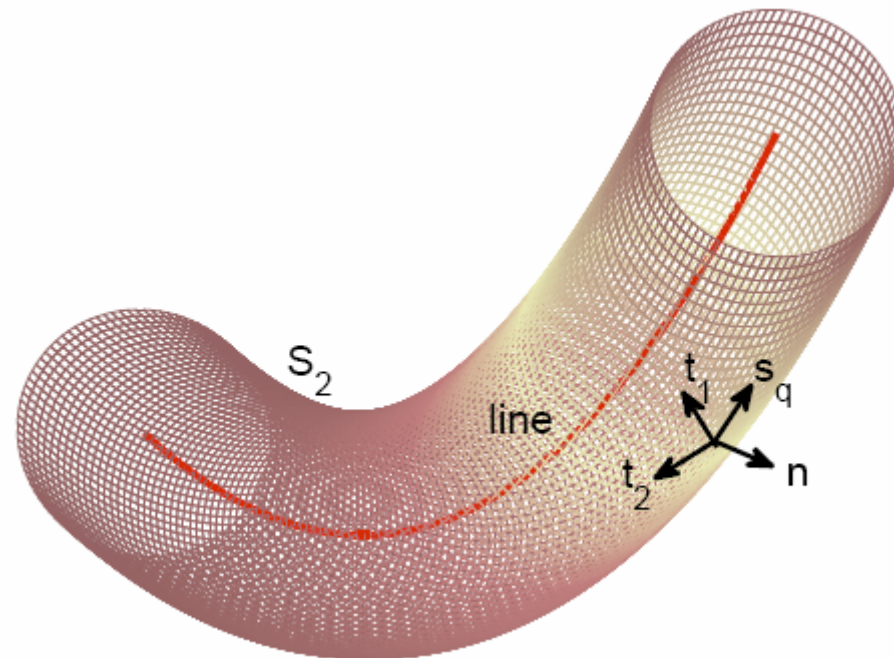
# Coordinate transformation



$$\hat{\mathbf{E}} = \Lambda^T \mathbf{E}, \quad \hat{\mathbf{H}} = \Lambda^T \mathbf{H},$$

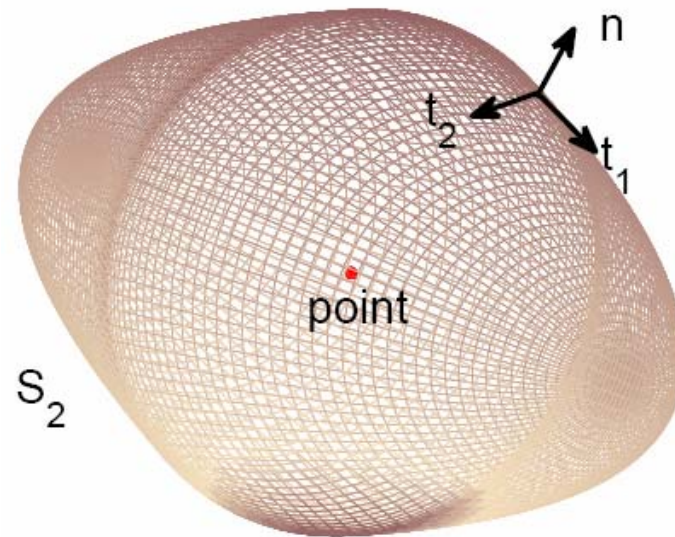


# Line-transformed cloak





# Point-transformed cloak





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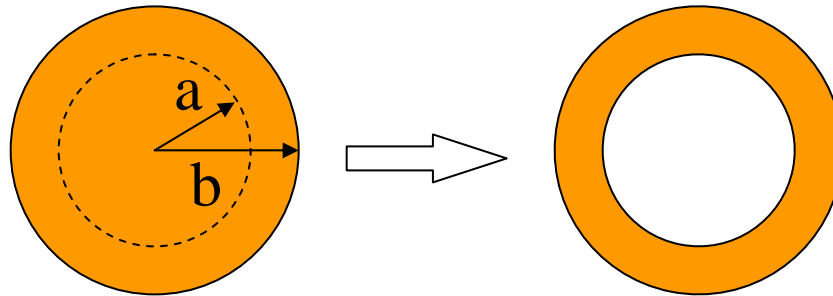
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# Ideal cylindrical cloak

## Principle:

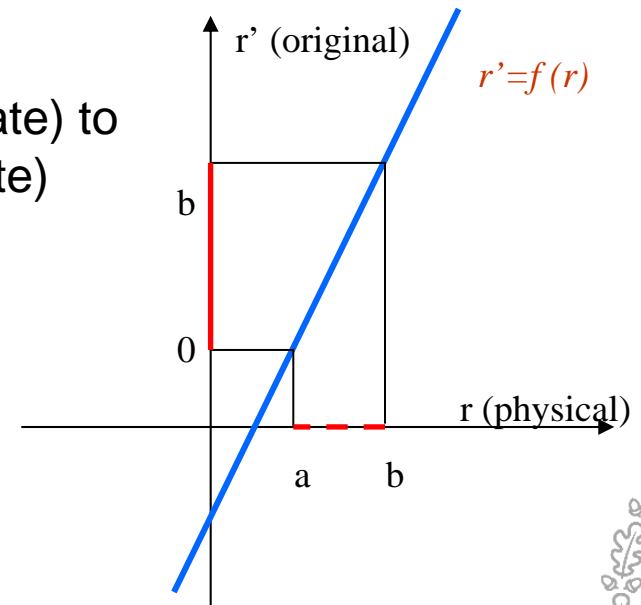
To compress EM fields within a cylindrical air region  $r' < b$  into the cylindrical annular region  $a < r < b$ .



## The simplest transformation:

from  $(r', \theta', z')$  (original cylindrical coordinate) to  $(r, \theta, z)$  (physical cylindrical coordinate)

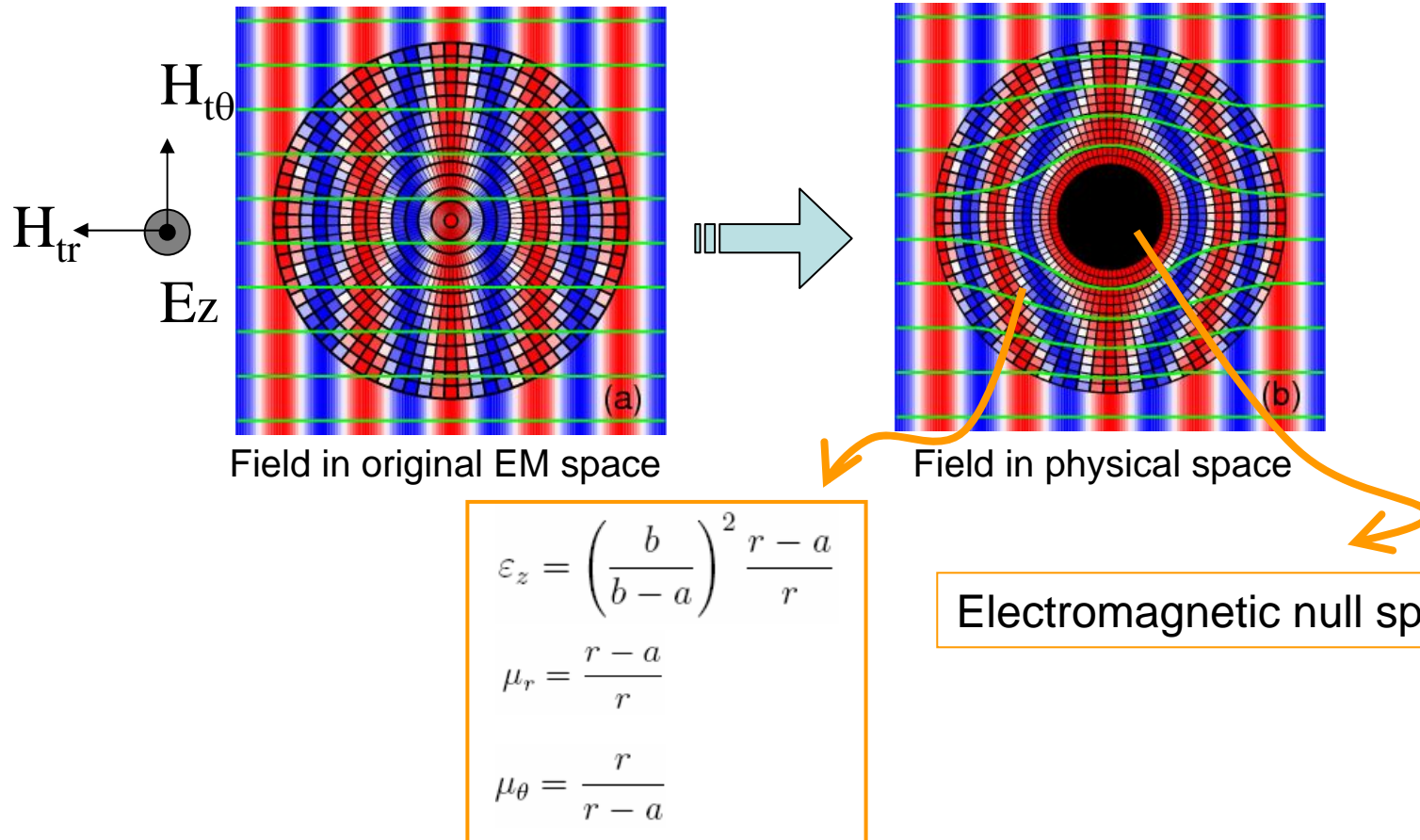
$$\begin{aligned} r &= \frac{b-a}{a} r' + a \\ \theta &= \theta' \\ z &= z' \end{aligned}$$





# Ideal cylindrical cloak

Colormap shows  $E_z$  field



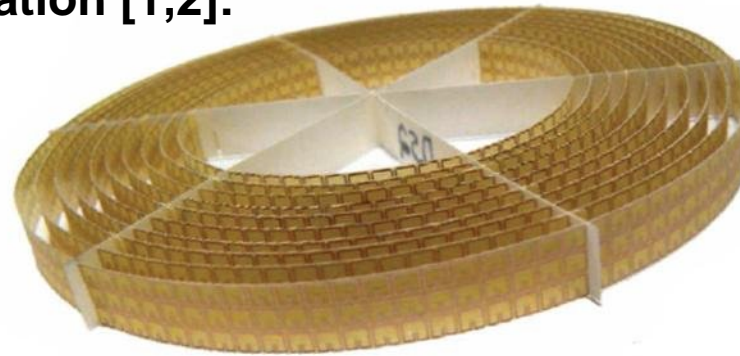
**Difficulty:**  $\epsilon_\theta$ ,  $\mu_\theta$ ,  $\epsilon_z$ , and  $\mu_z$  diverge at  $r=a$ !



# Simplified cylindrical cloaks

## Reasons for material simplification [1,2]:

1. To avoid infinite parameters
2. To alleviate metamaterial engineering task



See Ref. 1

The material parameters are simplified as

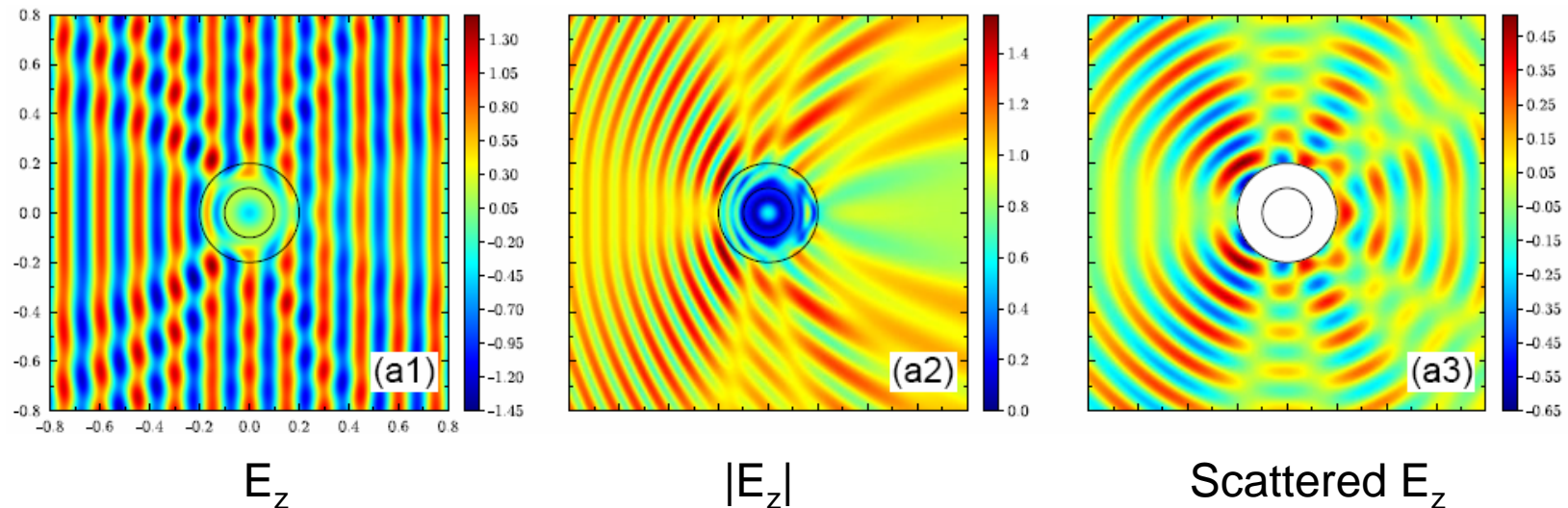
$$\begin{array}{l} \varepsilon_r = \mu_r = \frac{r-a}{r}, \\ \varepsilon_\theta = \mu_\theta = \frac{r}{r-a}, \\ \varepsilon_z = \mu_z = \left(\frac{b}{b-a}\right)^2 \frac{r-a}{r}. \end{array} \quad \Rightarrow \quad \begin{array}{l} \varepsilon_r = \mu_r = \left(\frac{r-a}{r}\right)^2, \\ \varepsilon_\theta = \mu_\theta = 1, \\ \varepsilon_z = \mu_z = \left(\frac{b}{b-a}\right)^2. \end{array}$$

$\mu_\theta \varepsilon_z$  and  $\mu_r \varepsilon_z$  invariant

1. D. Schurig, et al., Science **314**, 977 (2006).
2. W. Cai et al., "Optical cloaking with metamaterials," Nat. Photonics **1**, 224–227 (2007).



# Simplified cylindrical cloaks: not perfect



## Origin of the problem:

Wave equation ( $r$ -dependent) in ideal cloak medium:

$$\frac{d}{dr} \left( \frac{r}{\mu_\theta} \frac{d\Psi}{dr} \right) + k_0^2 r \epsilon_z \Psi - m^2 \frac{1}{r \mu_r} \Psi = 0.$$

Wave equation in simplified cloak medium:

$$\frac{1}{\mu_\theta \epsilon_z} \frac{d}{dr} \left( r \frac{d\Psi}{dr} \right) + k_0^2 r \Psi - m^2 \frac{1}{r \mu_r \epsilon_z} \Psi = 0.$$

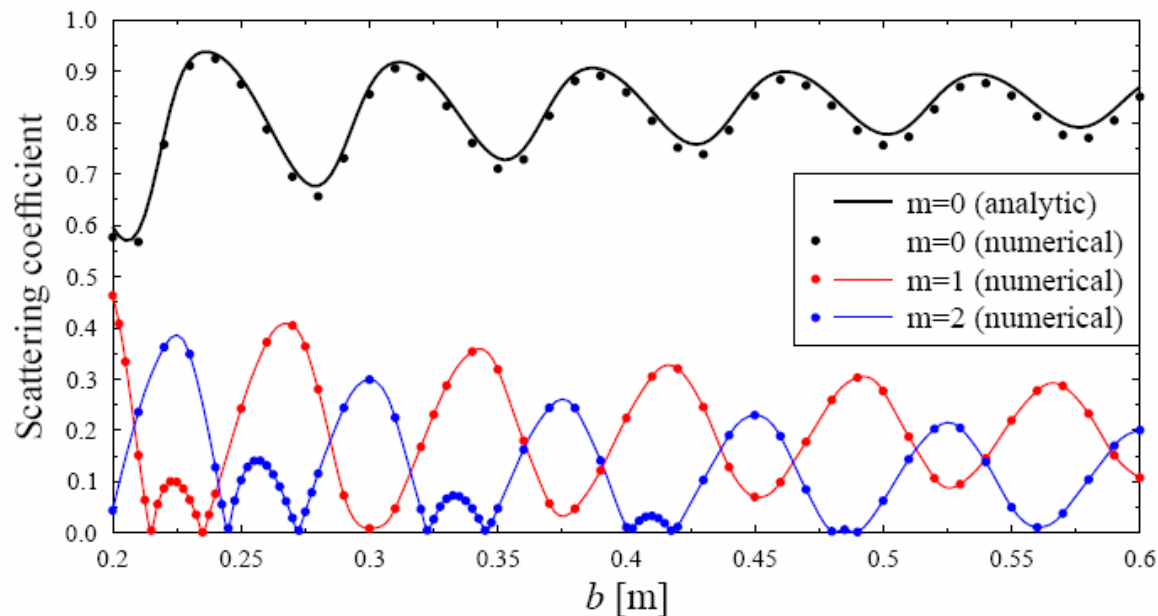
Wrong assumption:  $\mu_\theta$  is constant

where  $m$  is angular mode number



# Scattering coefficients

Scattering coefficients in different cylindrical orders w.r.t. *Outer radius*  $b$  [1]:



**Parameters:**  $f=2\text{GHz}$ ;  $a=0.1\text{m}$

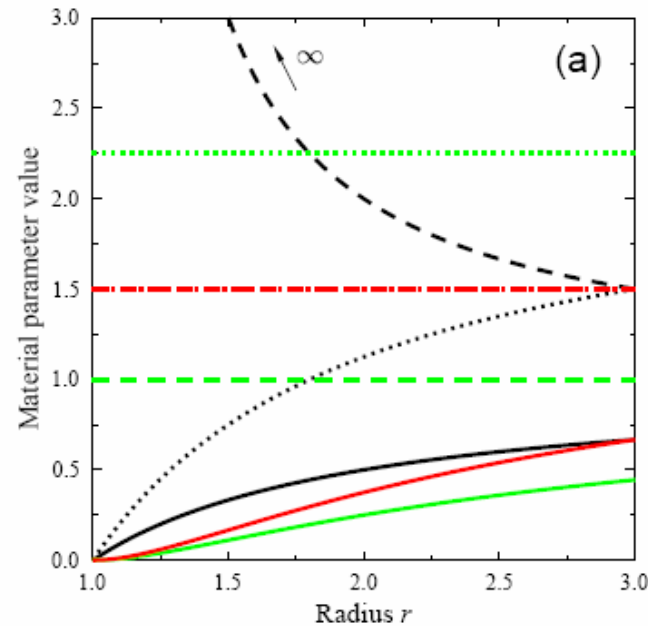


# Improved simplification

Ideal	Simplified [2]	Simplified (current work)
$\epsilon_r = \mu_r = \frac{r-a}{r}$	$\epsilon_r = \mu_r = \left(\frac{r-a}{r}\right)^2$	$\epsilon_r = \mu_r = \left(\frac{r-a}{r}\right)^2 \frac{b}{b-a}$
$\epsilon_\theta = \mu_\theta = \frac{r}{r-a}$	$\epsilon_\theta = \mu_\theta = 1$	$\epsilon_\theta = \mu_\theta = \frac{b}{b-a}$
$\epsilon_z = \mu_z = \left(\frac{b}{b-a}\right)^2 \frac{r-a}{r}$	$\epsilon_z = \mu_z = \left(\frac{b}{b-a}\right)^2$	$\epsilon_z = \mu_z = \frac{b}{b-a}$

Same condition:  $\mu_\theta \epsilon_z$  and  $\mu_r \epsilon_z$  invariant

—  $\mu_r$   
 - -  $\mu_\theta$   
 .....  $\epsilon_z$



## Advantage:

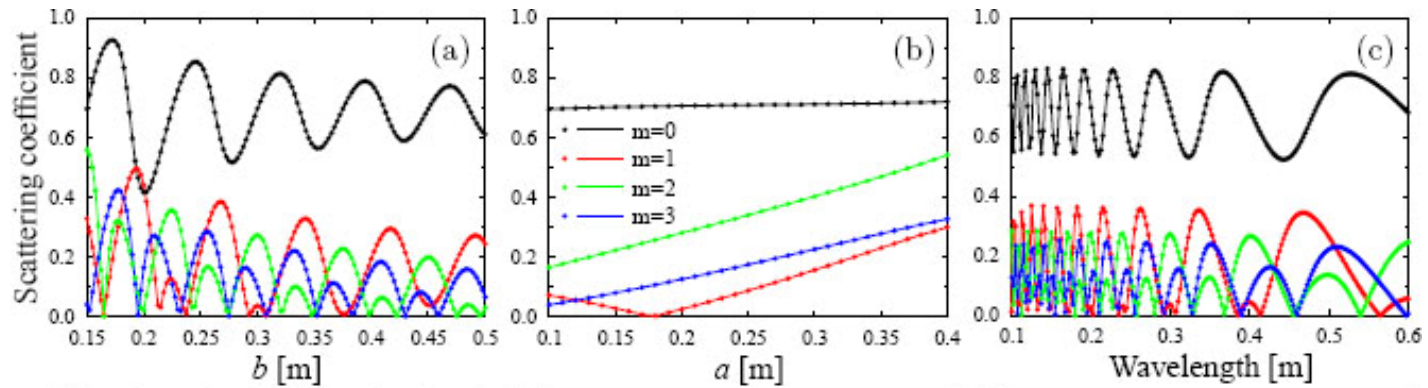
The outer interface is perfectly matched to exterior!

1. M. Yan, Z. Ruan, and M. Qiu, Opt. Express **15**, p. 17772 (2007).
2. D. Schurig, et al., Science **314**, 977 (2006).

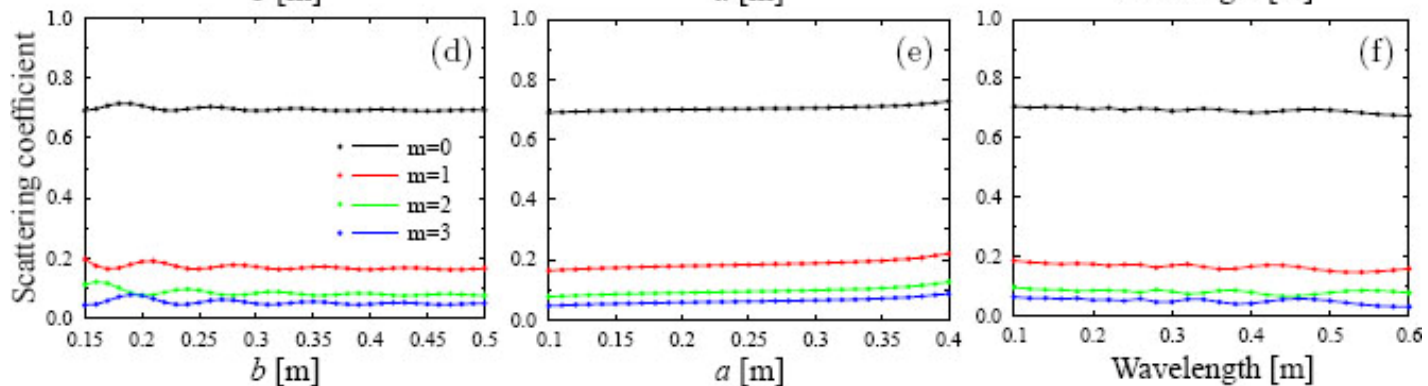


# Improved simplification

Previous  
simplified

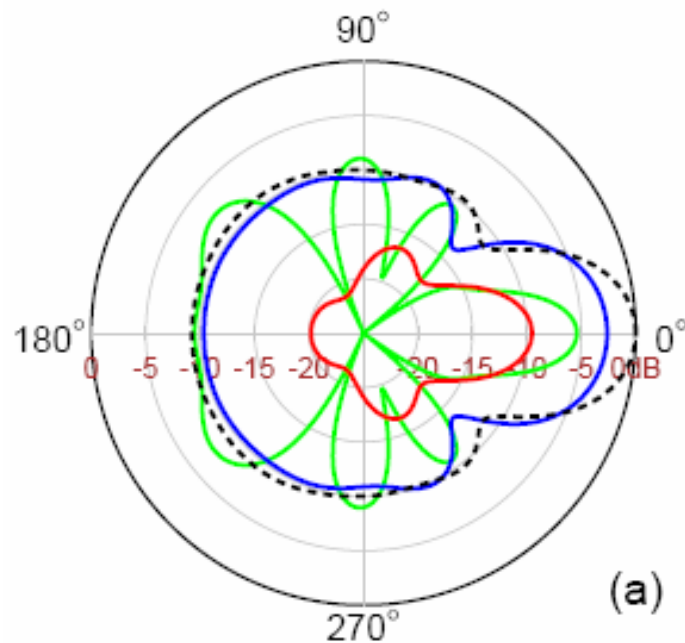


Improved





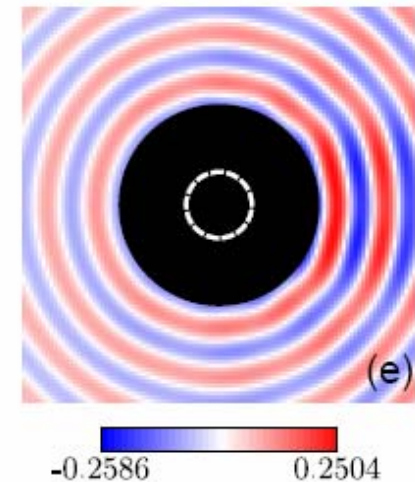
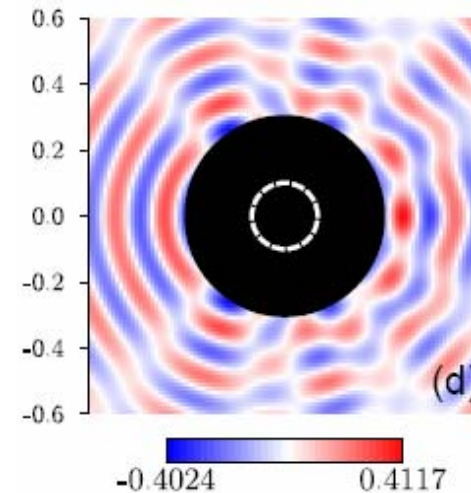
# Improved simplification



- ..... Bare cylinder
- Simplified linear cloak [2]
- Simplified quadratic cloak [3]
- Simplified linear cloak (improved)

Linear simplified cloak

Current work

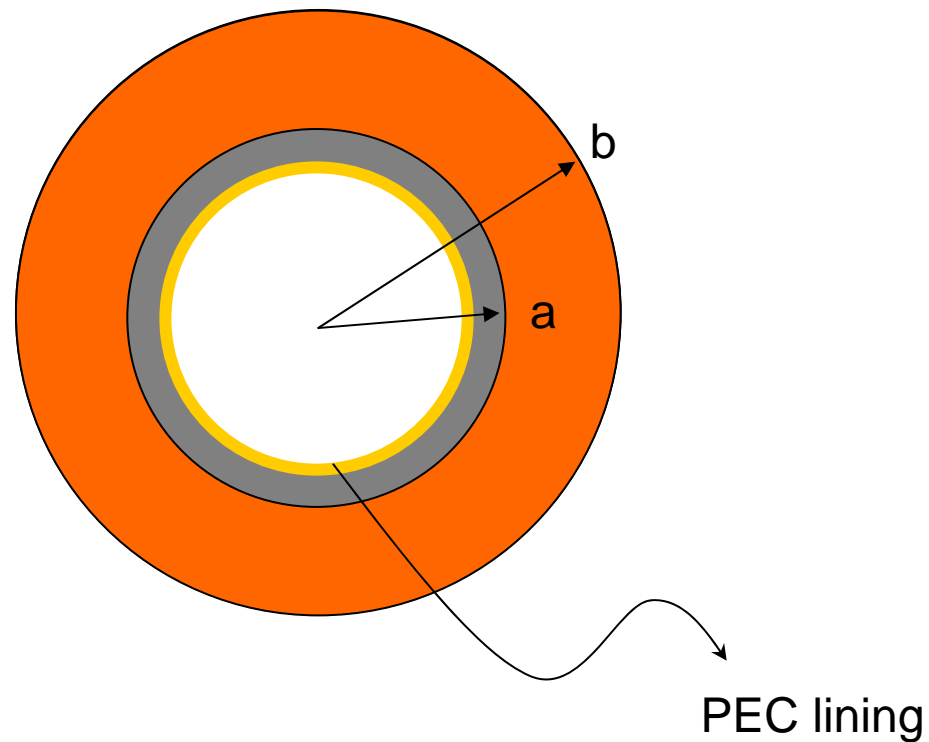


Scattered  $E_z$  field

1. M. Yan, Z. Ruan, and M. Qiu, Opt. Express **15**, p. 17772 (2007).
2. D. Schurig, et al., Science **314**, 977 (2006).
3. W. Cai et al., Appl. Phys. Lett. **91**, 111,105 (2007).



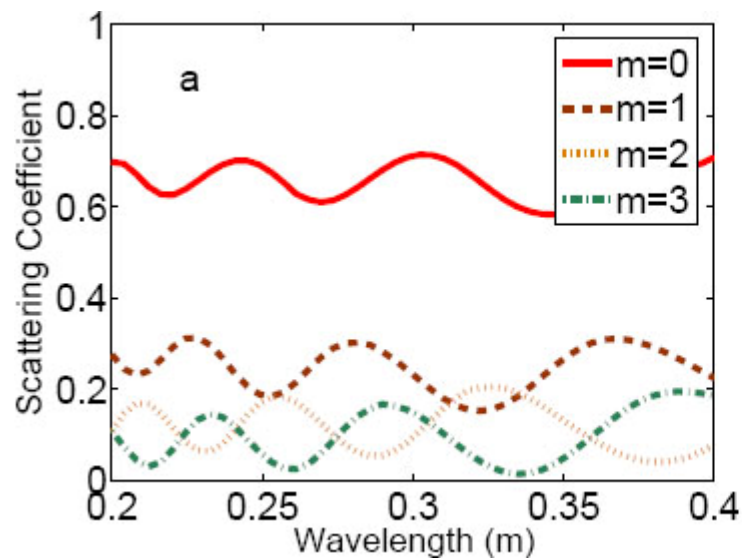
# Can we cancel the zeroth order scattering?



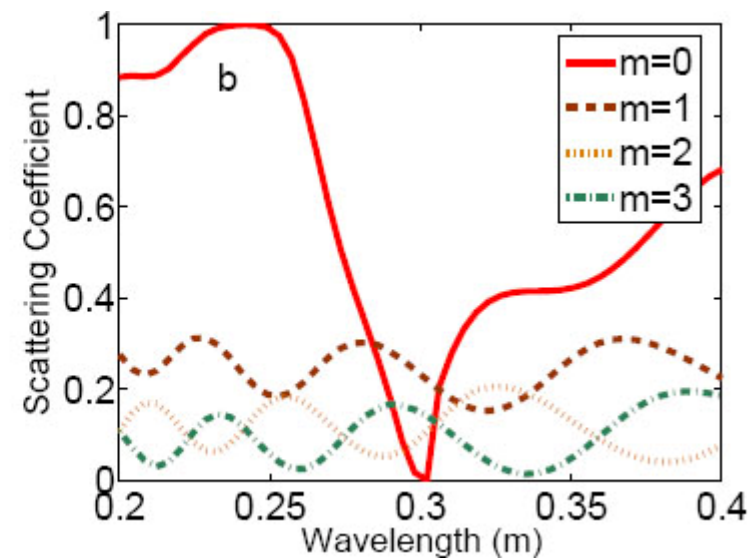


# Cancellation of zeroth-order scattering

without monopole cancellation



*with* monopole cancellation

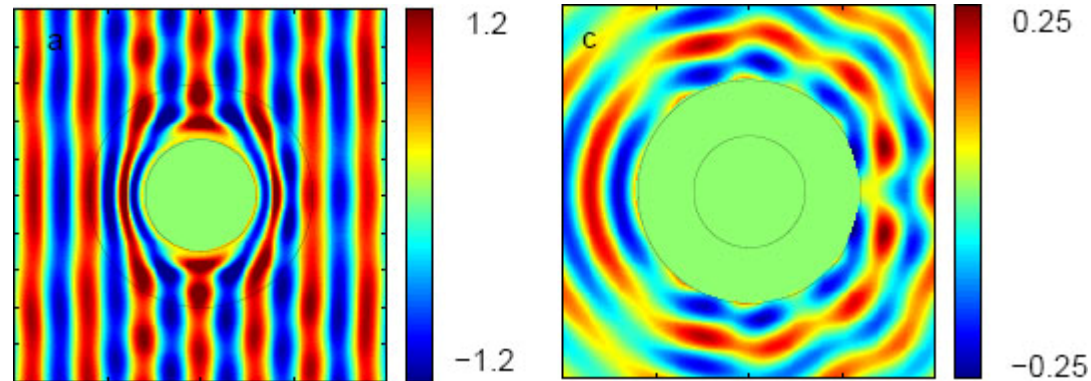


$a = 0.3\text{m}$ ,  $b = 0.6\text{m}$

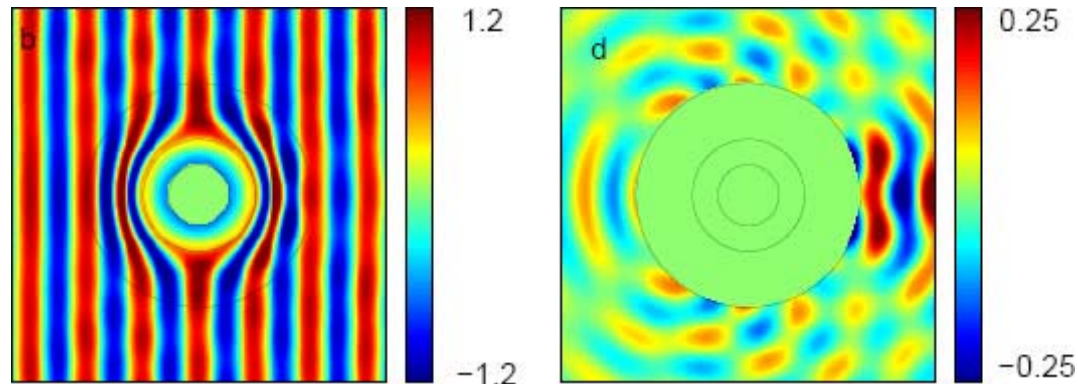


# Cancellation of zeroth-order scattering

Without  
cancellation



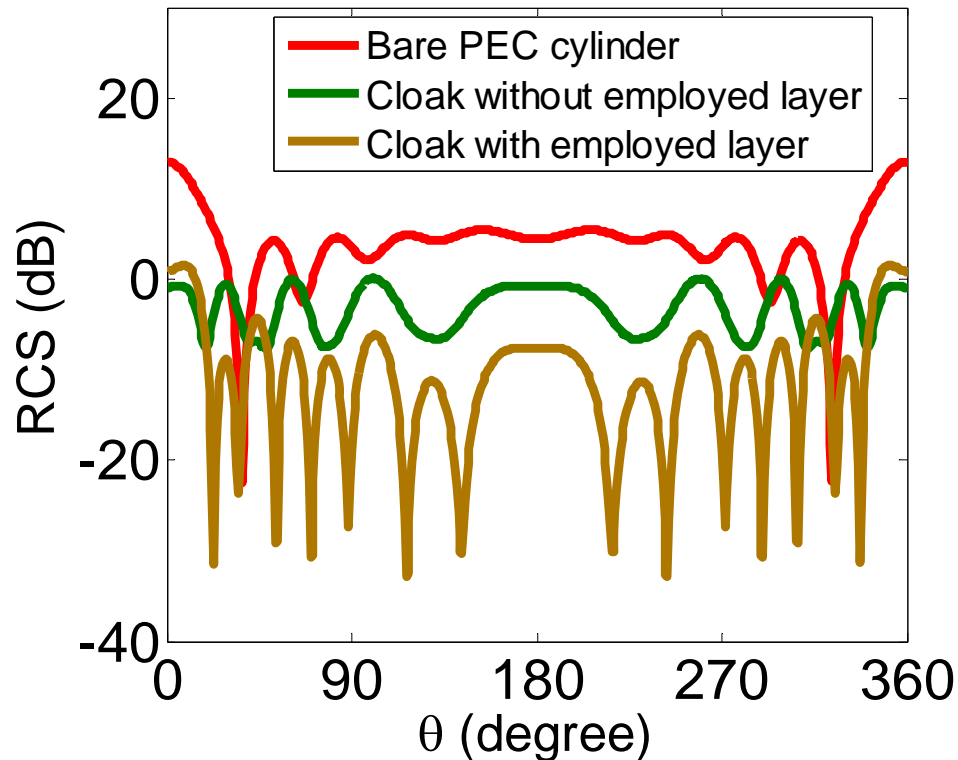
With  
cancellation



$$\lambda=0.3\text{m}$$



# Far field radiation



## Remarks:

- Mostly high order scatterings remain
- Wavelength dependent

## Ref:

W. Yan, M. Yan, and Min Qiu, "Non-magnetic simplified cylindrical cloak with near perfect invisibility by suppressing zeroth order scattering," Appl. Phys. Lett., in review.



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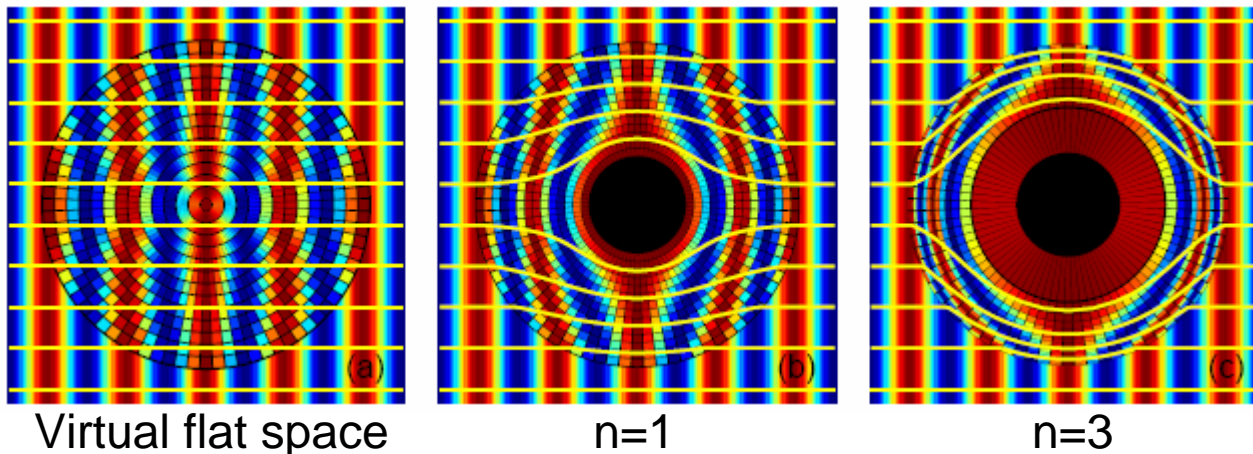


# High-order cylindrical invisibility cloak

One possible class of transformation may take the form of

$$r' = \frac{b}{(b-a)^n} (r-a)^n,$$

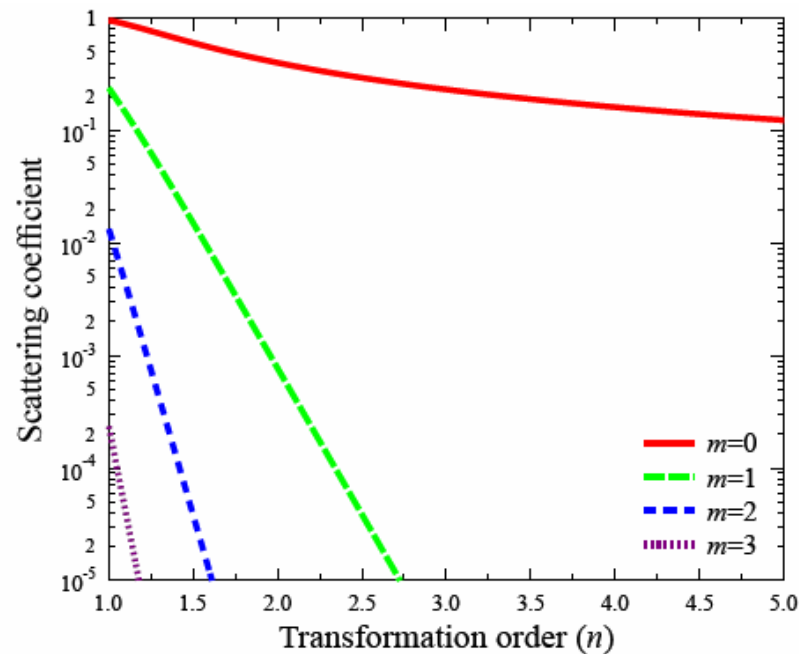
$n$ : transformation order.



$$\epsilon_r = \mu_r = \frac{r-a}{nr}, \quad \epsilon_\theta = \mu_\theta = \frac{nr}{r-a}, \quad \epsilon_z = \mu_z = \frac{nb^2(r-a)^{2n-1}}{(b-a)^{2n}r}.$$

Infinite parameters at  $r=a$ !

# Effect of transformation order on scattering coefficients



$a = 0.1\text{m}$

$b = 0.3\text{m}$

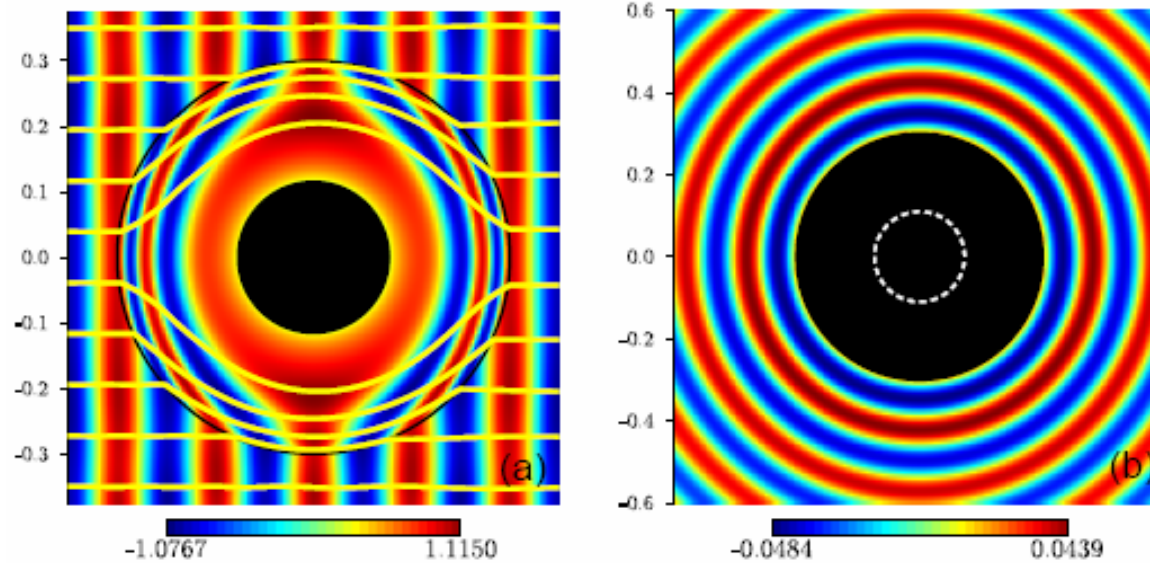
$\lambda = 0.15\text{m}$

Thickness of the layer peeled away:  $d=0.01\text{m}$

PEC lining is present



# Field distributions



$E_z$

Scattered  $E_z$

Transformation order used:  $n=3$

## Ref

M. Yan, W. Yan, L. Zhang, M. Qiu, Cylindrical Invisibility Cloak: Properties and Strategies for Practical Realization, NATO Advanced Research Workshop, Metamaterials for Secure Information and Communication Technologies, 7-10 May, 2008, Marrakech – Morocco.

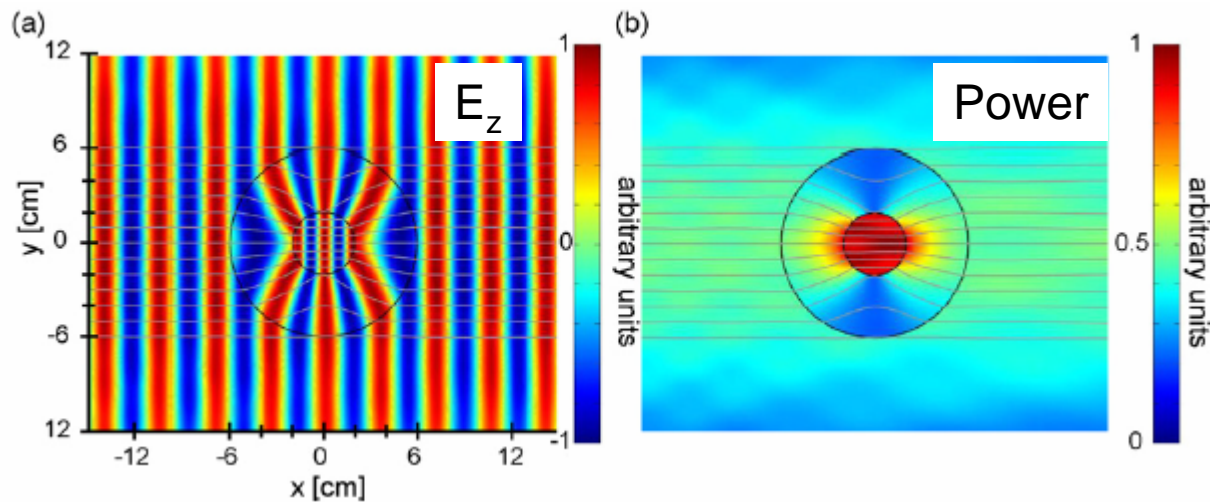
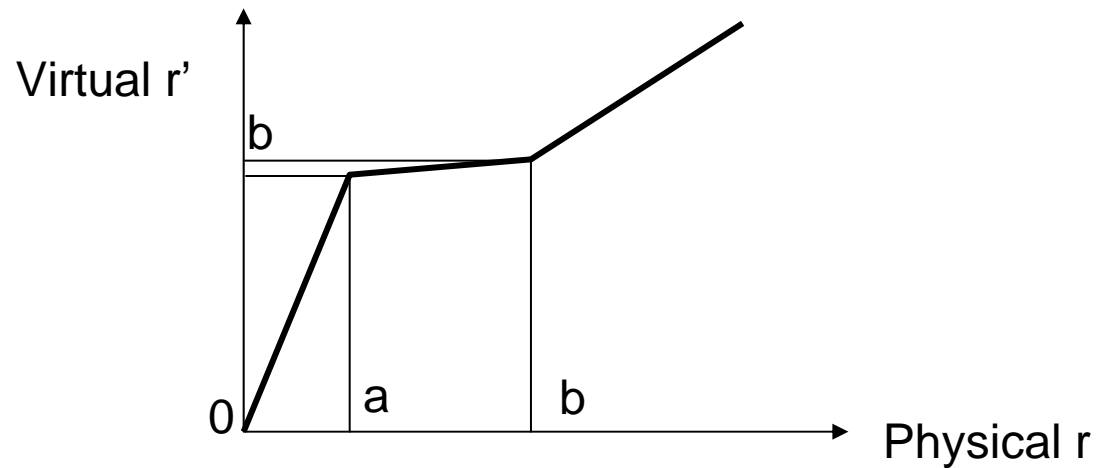


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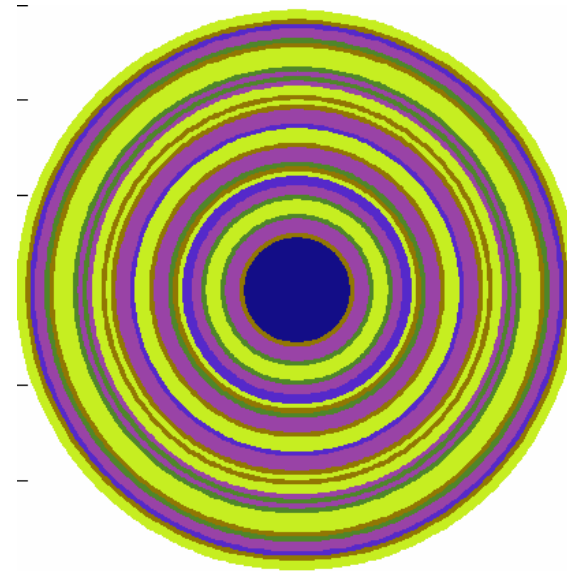
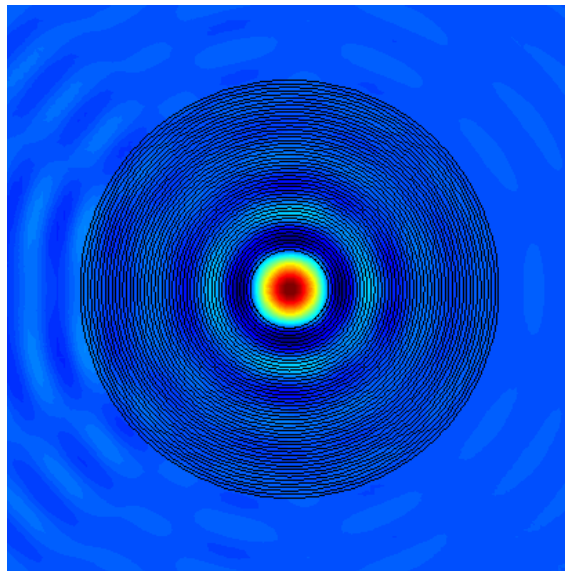
# EM concentrator



- Application:**
- Photovoltaic fiber
  - Fluid heating



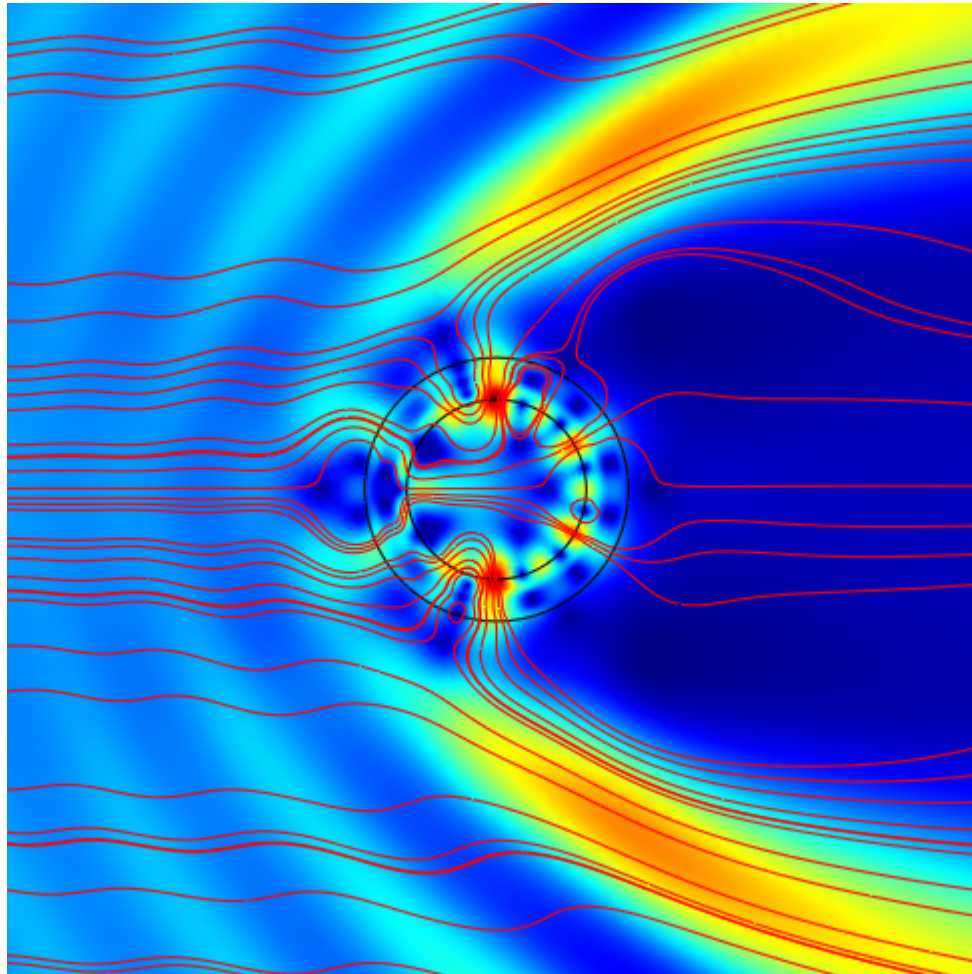
# Resonance based EM concentrator



- Enhancement factor 230
- Wavelength dependant



# Resonance based EM concentrator



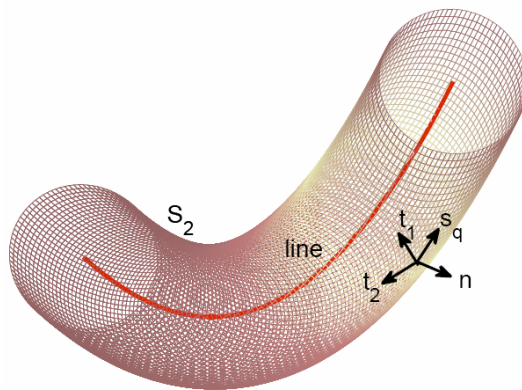


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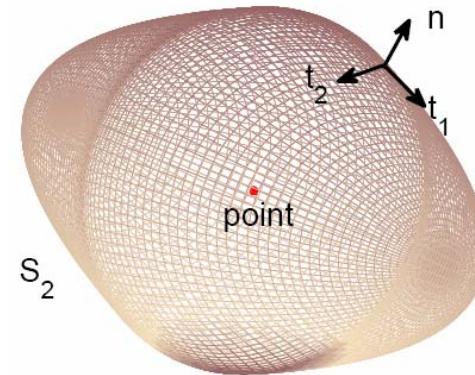
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# Arbitrarily-shaped cloaks

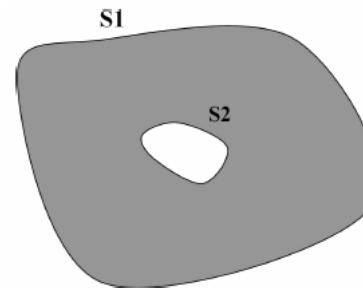


Blowing-up a line



Blowing-up a point

General cross-section:



**Ref:**

Wei Yan, Min Yan, Zhichao Ruan, Min Qiu, "Coordinate transformation makes perfect invisibility cloak with arbitrary shape," New Journal of Physics, vol. 10, 043040 (2008).



# Arbitrarily-shaped cloak

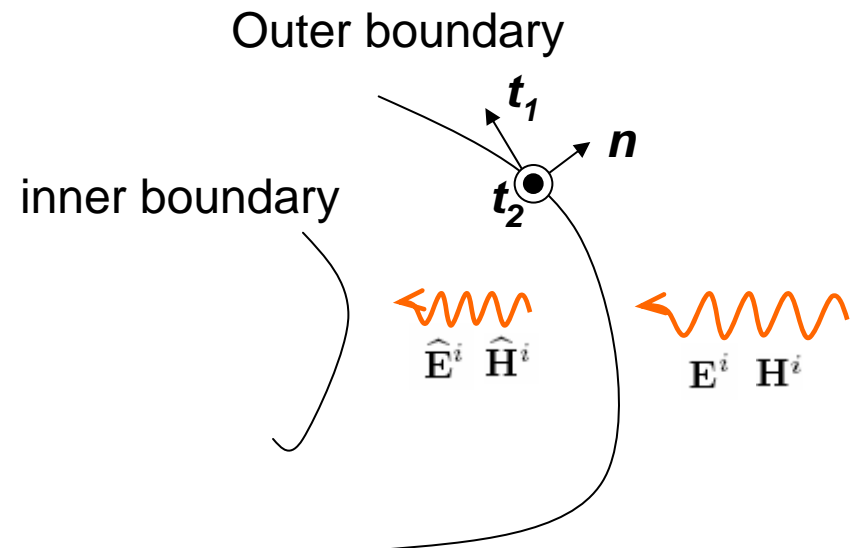
For both line-transformed or point-transformed cloaks,

$$\hat{\mathbf{E}}^i = \Lambda^T \mathbf{E}^i, \quad \hat{\mathbf{H}}^i = \Lambda^T \mathbf{H}^i,$$

At outer boundary:

$$\hat{E}_{t_1}^i = \mathbf{E}^i \cdot \hat{t}_1, \quad \hat{H}_{t_1}^i = \mathbf{H}^i \cdot \hat{t}_1,$$

$$\hat{E}_{t_2}^i = \mathbf{E}^i \cdot \hat{t}_2, \quad \hat{H}_{t_2}^i = \mathbf{H}^i \cdot \hat{t}_2,$$





# Arbitrarily-shaped cloak

For line-transformed cloak,

$$\hat{E}_{t1}^i = \hat{H}_{t1}^i = 0.$$

However, the other components of fields are not zero. In particular,

$$\hat{E}_{t2}^i = (\hat{s} \cdot \hat{t}_2)[B_1, B_2, B_3]\mathbf{E}^i,$$

$$\hat{H}_{t2}^i = (\hat{s} \cdot \hat{t}_2)[B_1, B_2, B_3]\mathbf{H}^i,$$

$$\hat{E}_n^i = [F_1 + B_1(\hat{s} \cdot \hat{n}), F_2 + B_2(\hat{s} \cdot \hat{n}), F_3 + B_3(\hat{s} \cdot \hat{n})]\mathbf{E}^i,$$

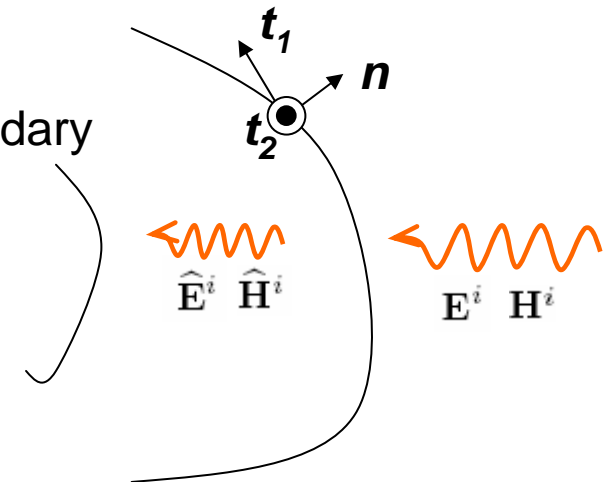
$$\hat{H}_n^i = [F_1 + B_1(\hat{s} \cdot \hat{n}), F_2 + B_2(\hat{s} \cdot \hat{n}), F_3 + B_3(\hat{s} \cdot \hat{n})]\mathbf{H}^i,$$

with

$$B_i = \sqrt{\partial b_i / \partial q_1^2 + \partial b_i / \partial q_2^2 + \partial b_i / \partial q_3^2}.$$

inner boundary

Outer boundary



**Remark:** Surface current will be induced at the inner surface.



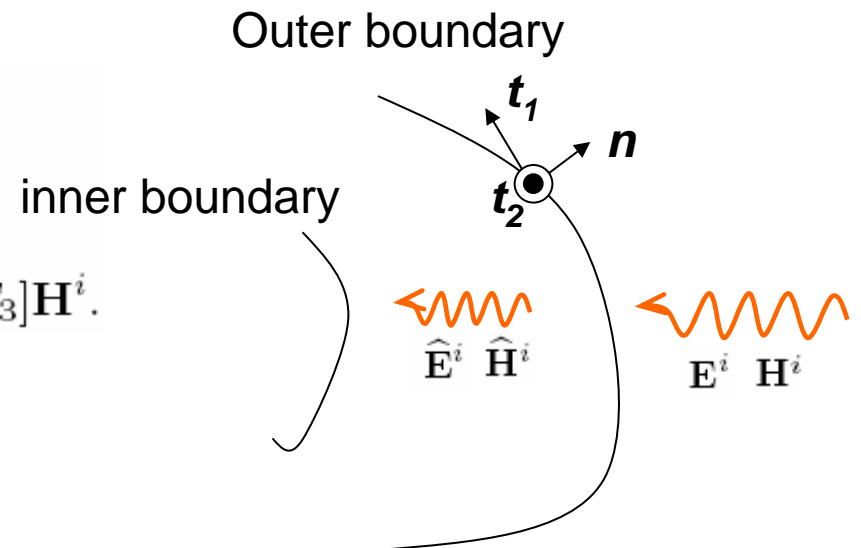
# Arbitrarily-shaped cloak

For point-transformed cloak,

$$\hat{E}_{t_1}^i = \hat{H}_{t_1}^i = 0,$$

$$\hat{E}_{t_2}^i = \hat{H}_{t_2}^i = 0,$$

$$\hat{E}_n^i = [F_1, F_2, F_3] \mathbf{E}^i, \quad \hat{H}_n^i = [F_1, F_2, F_3] \mathbf{H}^i.$$



All tangential fields are zero.

Therefore no field discontinuity exists for perfect cloaking.



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# Project plan

We are here!



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Project report								



# Outcome

## **Papers:**

1. Wei Yan, Min Yan, Zhichao Ruan, Min Qiu, "Coordinate transformation makes perfect invisibility cloak with arbitrary shape," New Journal of Physics, vol. 10, 043040 (2008).
2. M. Yan, W. Yan, L. Zhang, M. Qiu, Cylindrical Invisibility Cloak: Properties and Strategies for Practical Realization, NATO Advanced Research Workshop, Metamaterials for Secure Information and Communication Technologies, 7-10 May, 2008, Marrakech, Morocco.
3. W. Yan, M. Yan, and Min Qiu, "Non-magnetic simplified cylindrical cloak with near perfect invisibility by suppressing zeroth order scattering," Appl. Phys. Lett., in review.

## **Code:**

Matlab code for scattering calculation of multilayered cylindrical structures.



Thank you!



Questions?

