



Manipulation of Lightwave Through Coordinate Transformation

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Ariadna bidder code: **21290**





Project plan

We are here!



Tasks	Time (mth)	0.5	1	1.5	2	2.5	3	3.5	4
Cylindrical cloak: Simplification									
Cylindrical cloak: Effect of order number									
<i>Electromagnetic Concentrator</i>									
Arbitrarily-shaped cloak									
Project report									



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0. Introduction	8
1. Cylindrical cloak: Simplification	15'
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4. Arbitrarily-shaped cloak	6'
5. Conclusion	4'
Total:	45'



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Introduction



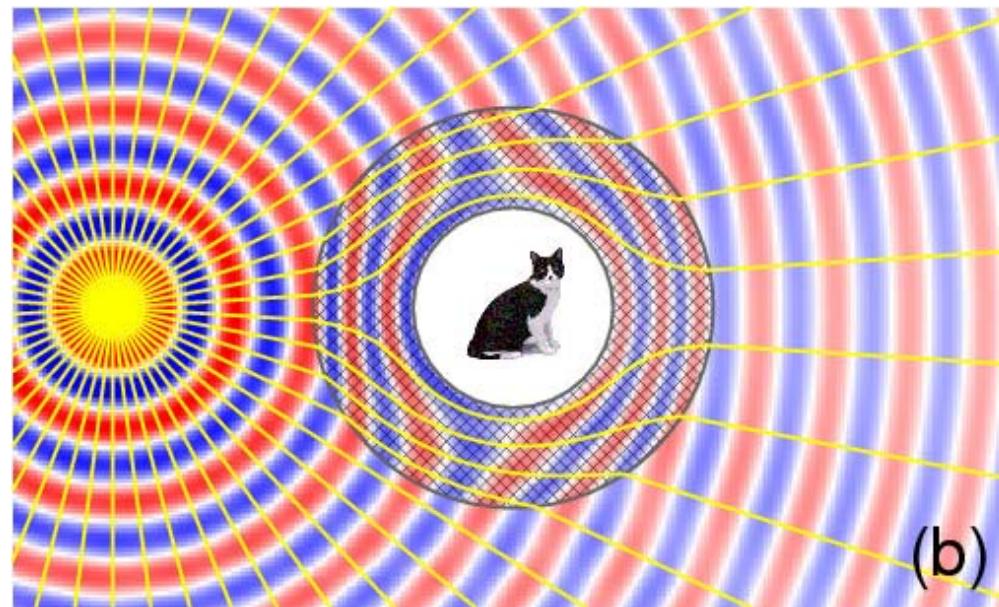
Scattering → visibility



? → invisibility



Coordinate transformation approach



(b)

1. J. B. Pendry, D. Schurig, and D. R. Smith, *Science* **312**, p.1780 (2006).
2. U. Leonhardt, *Science* **312**, p. 1777 (2006).



Coordinate transformation

Original Cartesian coordinate (x, y, z)

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}, \quad \nabla \times \mathbf{H} = \frac{\partial \mathbf{D}}{\partial t} + \mathbf{j}, \quad \nabla \cdot \mathbf{D} = \rho, \quad \nabla \cdot \mathbf{B} = 0.$$

$$\mathbf{D} = \epsilon_0 \bar{\bar{\epsilon}} \cdot \mathbf{E}, \quad \mathbf{B} = \mu_0 \bar{\bar{\mu}} \cdot \mathbf{H},$$

Jacobian transformation matrix

$$\Lambda = \begin{bmatrix} \frac{\partial x}{\partial q_1} & \frac{\partial x}{\partial q_2} & \frac{\partial x}{\partial q_3} \\ \frac{\partial y}{\partial q_1} & \frac{\partial y}{\partial q_2} & \frac{\partial y}{\partial q_3} \\ \frac{\partial z}{\partial q_1} & \frac{\partial z}{\partial q_2} & \frac{\partial z}{\partial q_3} \end{bmatrix}.$$

New coordinate (q_1, q_2, q_3)

$$x = f_1(q_1, q_2, q_3), \quad y = f_2(q_1, q_2, q_3), \quad z = f_3(q_1, q_2, q_3).$$

The Maxwell equations can take the invariant form as

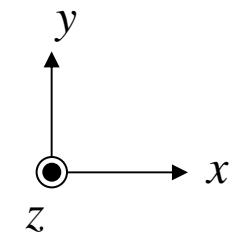
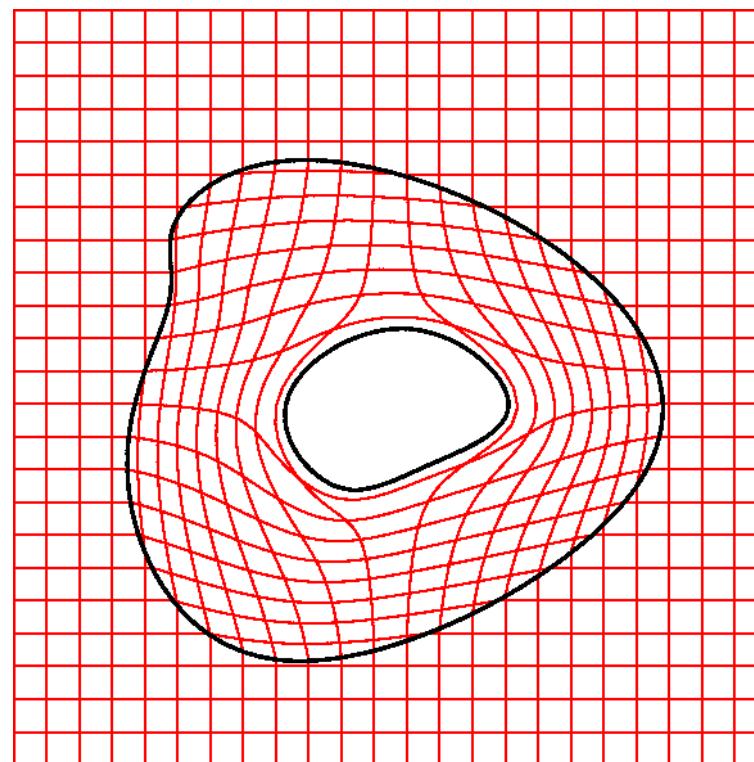
$$\nabla_q \times \hat{\mathbf{E}} = -\frac{\partial \hat{\mathbf{B}}}{\partial t}, \quad \nabla_q \times \hat{\mathbf{H}} = \frac{\partial \hat{\mathbf{D}}}{\partial t} + \hat{\mathbf{j}}, \quad \nabla_q \cdot \hat{\mathbf{D}} = \hat{\rho}, \quad \nabla_q \cdot \hat{\mathbf{B}} = 0$$

with

$$\hat{\bar{\epsilon}} = \det(\Lambda)(\Lambda)^{-1} \bar{\bar{\epsilon}} \Lambda^{-T}, \quad \hat{\bar{\mu}} = \det(\Lambda)(\Lambda)^{-1} \bar{\bar{\mu}} \Lambda^{-T},$$

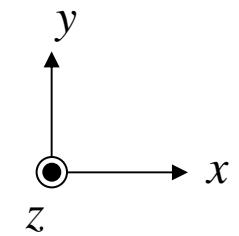
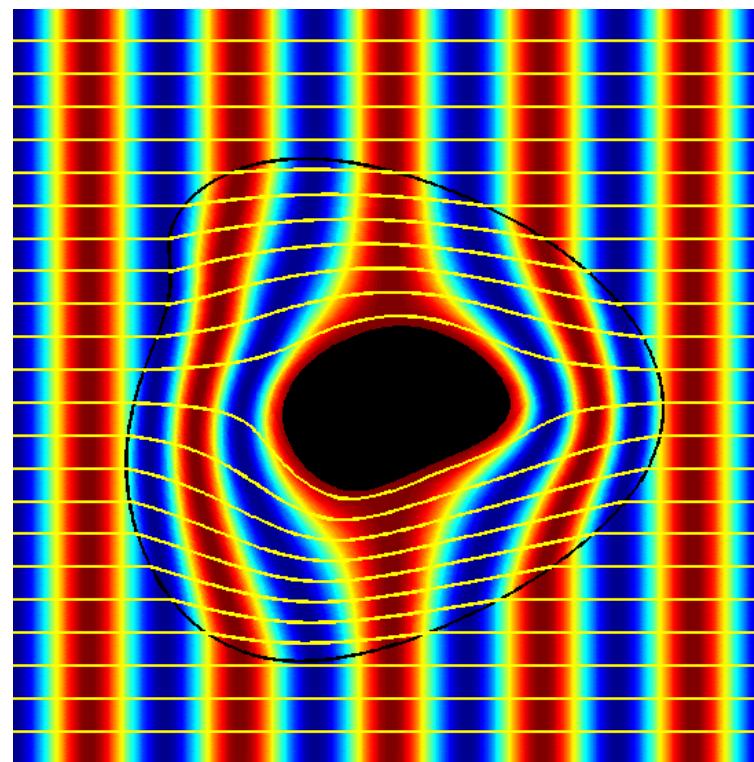


Coordinate transformation





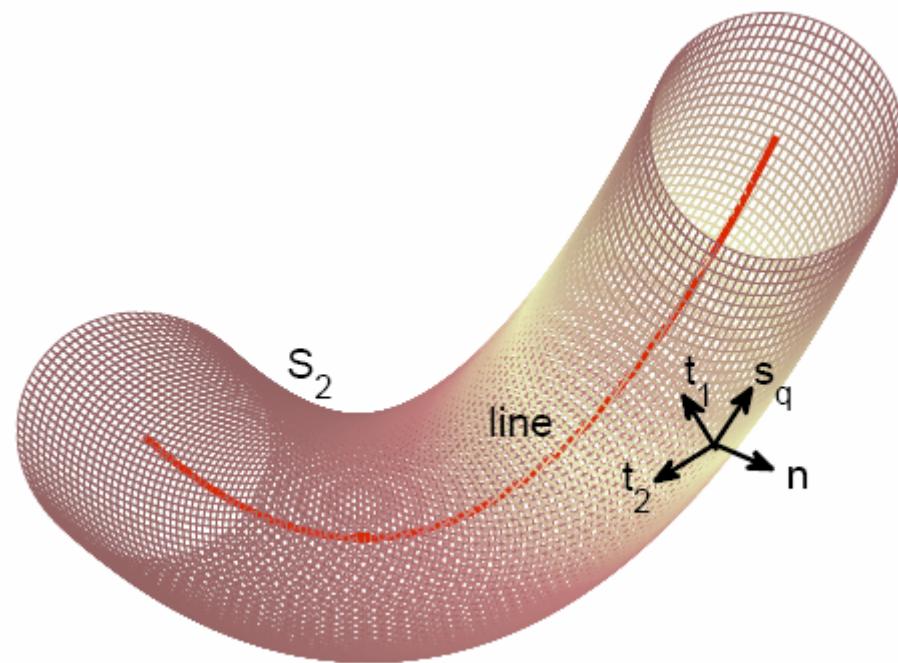
Coordinate transformation



$$\hat{\mathbf{E}} = \Lambda^T \mathbf{E}, \quad \hat{\mathbf{H}} = \Lambda^T \mathbf{H},$$



Line-transformed cloak





Point-transformed cloak

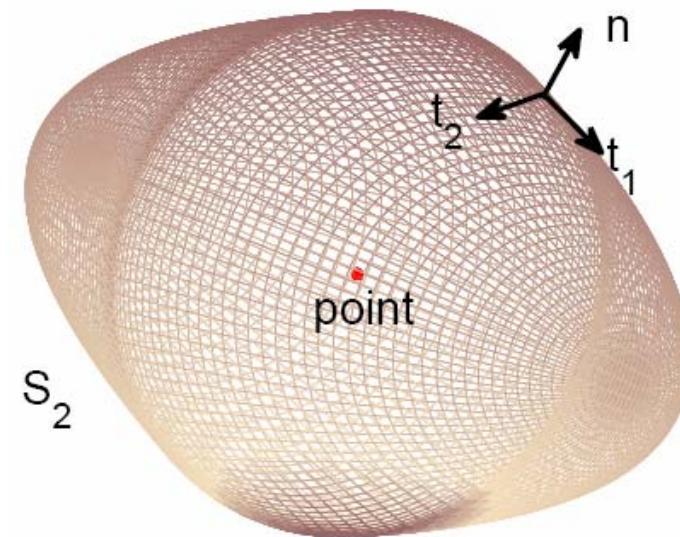




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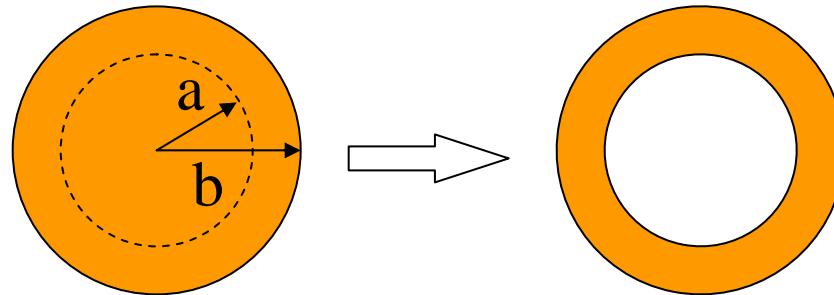
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Ideal cylindrical cloak

Principle:

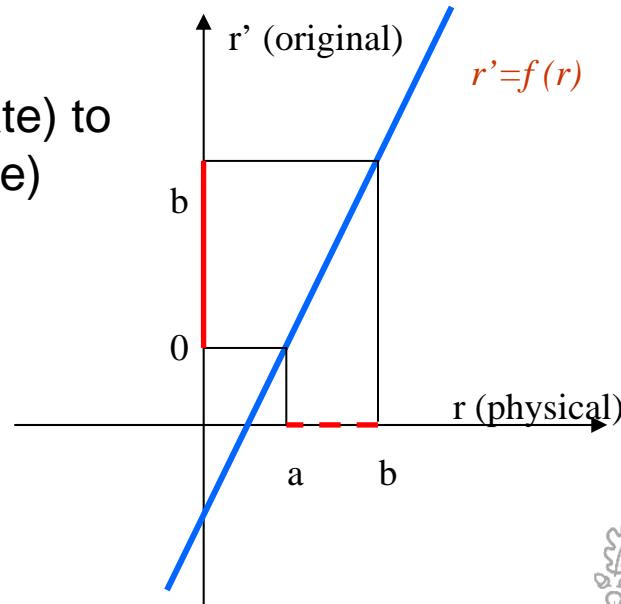
To compress EM fields within a cylindrical air region $r' < b$ into the cylindrical annular region $a < r < b$.



The simplest transformation:

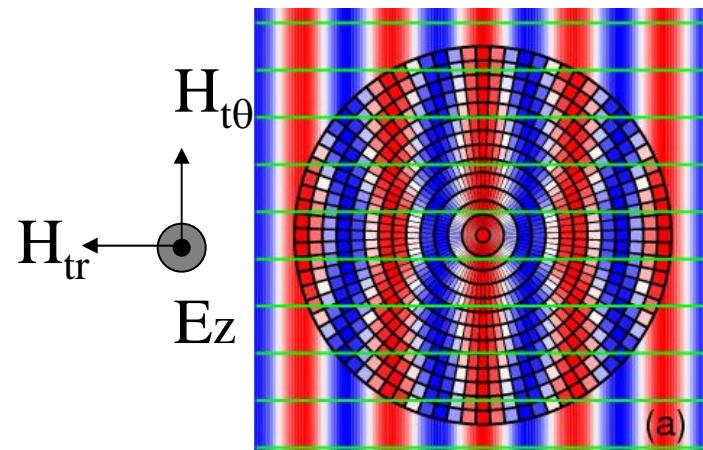
from (r', θ', z') (original cylindrical coordinate) to
 (r, θ, z) (physical cylindrical coordinate)

$$\begin{aligned} r &= \frac{b-a}{a} r' + a \\ \theta &= \theta' \\ z &= z' \end{aligned}$$

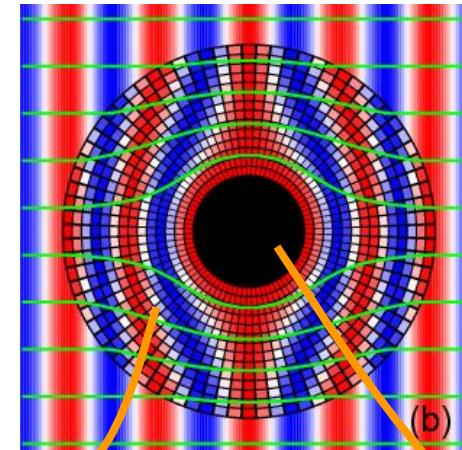
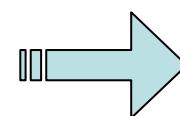




Ideal cylindrical cloak



Field in original EM space



Field in physical space

Colormap shows E_z field

$$\epsilon_z = \left(\frac{b}{b-a} \right)^2 \frac{r-a}{r}$$

$$\mu_r = \frac{r-a}{r}$$

$$\mu_\theta = \frac{r}{r-a}$$

Electromagnetic null space!

Difficulty: ϵ_θ , μ_θ , ϵ_z , and μ_z diverge at $r=a$!



Simplified cylindrical cloaks

Reasons for material simplification [1,2]:

1. To avoid infinite parameters
2. To alleviate metamaterial engineering task



See Ref. 1

The material parameters are simplified as

$$\begin{aligned}\varepsilon_r &= \mu_r = \frac{r-a}{r}, \\ \varepsilon_\theta &= \mu_\theta = \frac{r}{r-a}, \\ \varepsilon_z &= \mu_z = \left(\frac{b}{b-a}\right)^2 \frac{r-a}{r}.\end{aligned}$$



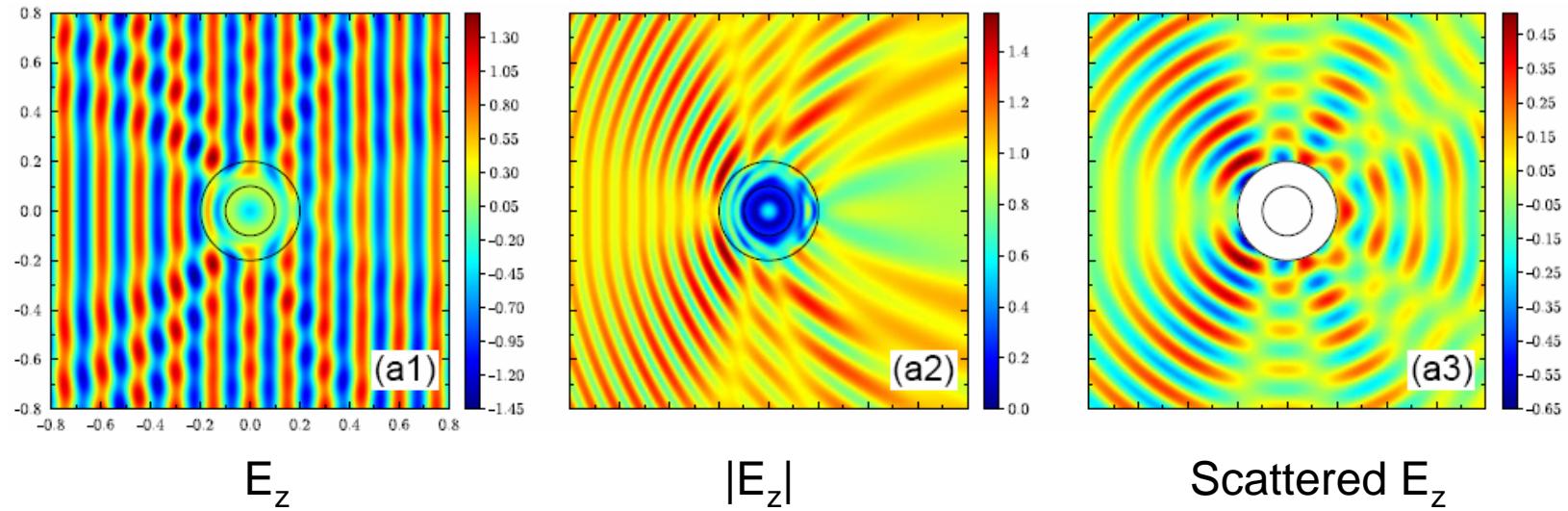
$$\begin{aligned}\varepsilon_r &= \mu_r = \left(\frac{r-a}{r}\right)^2, \\ \varepsilon_\theta &= \mu_\theta = 1, \\ \varepsilon_z &= \mu_z = \left(\frac{b}{b-a}\right)^2.\end{aligned}$$

$\mu_\theta \varepsilon_z$ and $\mu_r \varepsilon_z$ invariant

1. D. Schurig, et al., Science **314**, 977 (2006).
2. W. Cai et al., "Optical cloaking with metamaterials," Nat. Photonics **1**, 224–227 (2007).



Simplified cylindrical cloaks: not perfect



Origin of the problem:

Wave equation (r -dependent) in ideal cloak medium:

$$\frac{d}{dr} \left(\frac{r}{\mu_\theta} \frac{d\Psi}{dr} \right) + k_0^2 r \epsilon_z \Psi - m^2 \frac{1}{r \mu_r} \Psi = 0.$$

Wave equation in simplified cloak medium:

$$\frac{1}{\mu_\theta \epsilon_z} \frac{d}{dr} \left(r \frac{d\Psi}{dr} \right) + k_0^2 r \Psi - m^2 \frac{1}{r \mu_r \epsilon_z} \Psi = 0.$$

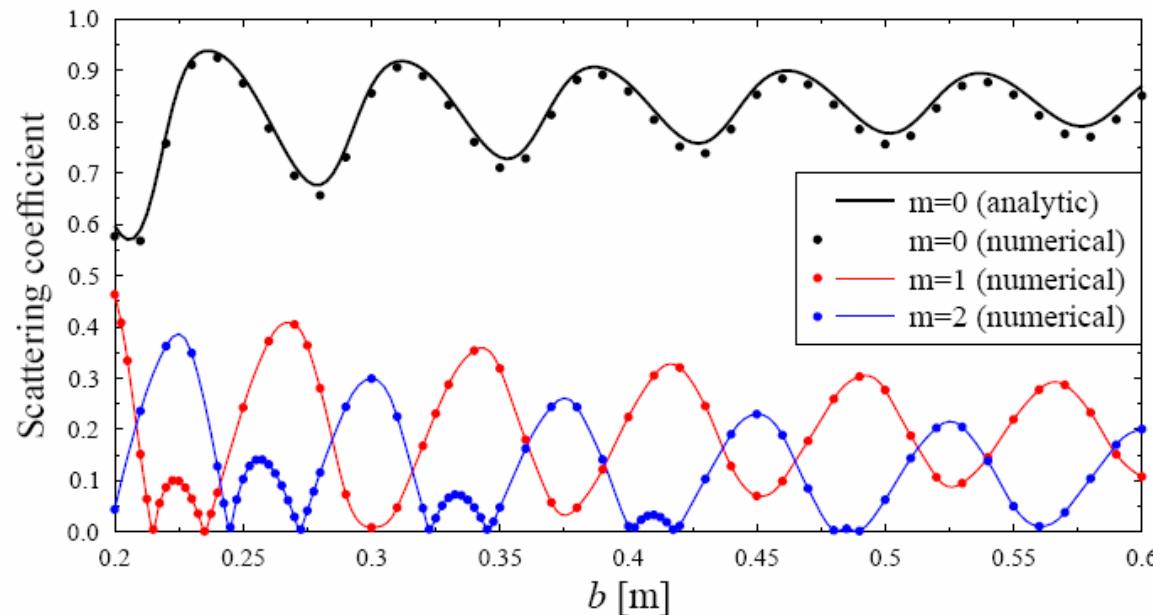
Wrong assumption: μ_θ is constant

where m is angular mode number



Scattering coefficients

Scattering coefficients in different cylindrical orders w.r.t. *Outer radius b* [1]:



Parameters: $f=2\text{GHz}$; $a=0.1\text{m}$

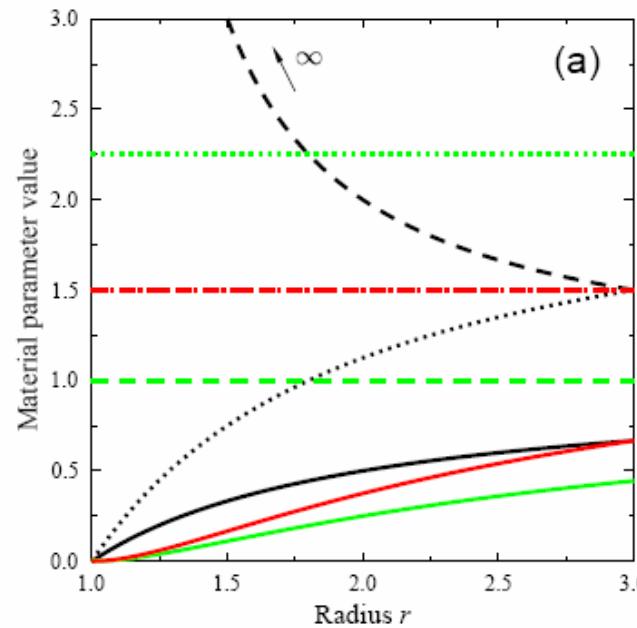


Improved simplification

Ideal	Simplified [2]	Simplified (current work)
$\epsilon_r = \mu_r = \frac{r-a}{r}$	$\epsilon_r = \mu_r = \left(\frac{r-a}{r}\right)^2$	$\epsilon_r = \mu_r = \left(\frac{r-a}{r}\right)^2 \frac{b}{b-a}$
$\epsilon_\theta = \mu_\theta = \frac{r}{r-a}$	$\epsilon_\theta = \mu_\theta = 1$	$\epsilon_\theta = \mu_\theta = \frac{b}{b-a}$
$\epsilon_z = \mu_z = \left(\frac{b}{b-a}\right)^2 \frac{r-a}{r}$	$\epsilon_z = \mu_z = \left(\frac{b}{b-a}\right)^2$	$\epsilon_z = \mu_z = \frac{b}{b-a}$

Same condition: $\mu_\theta \epsilon_z$ and $\mu_r \epsilon_z$ invariant

— μ_r
- - - μ_θ
- · - · - ϵ_z



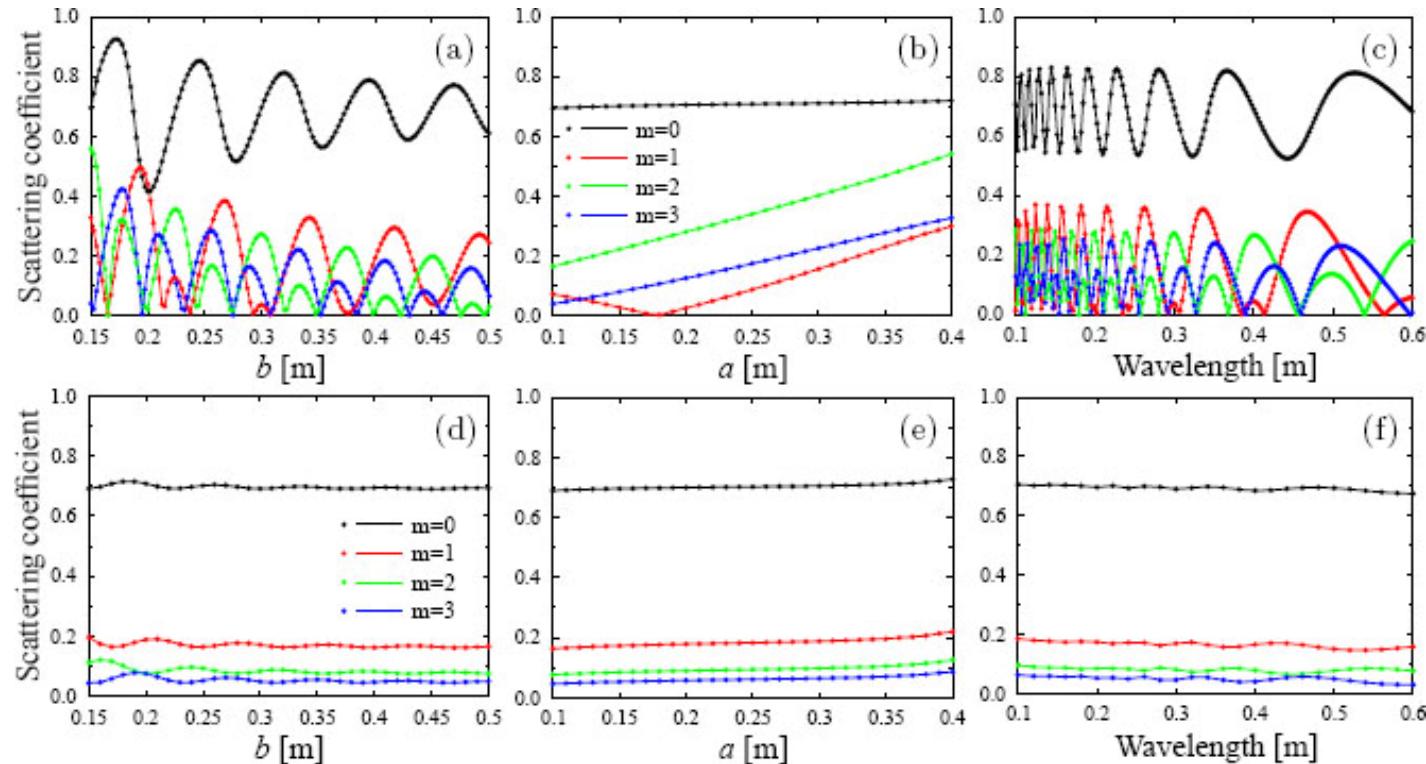
Advantage:
The outer interface is perfectly matched to exterior!

1. M. Yan, Z. Ruan, and M. Qiu, Opt. Express **15**, p. 17772 (2007).
2. D. Schurig, et al., Science **314**, 977 (2006).



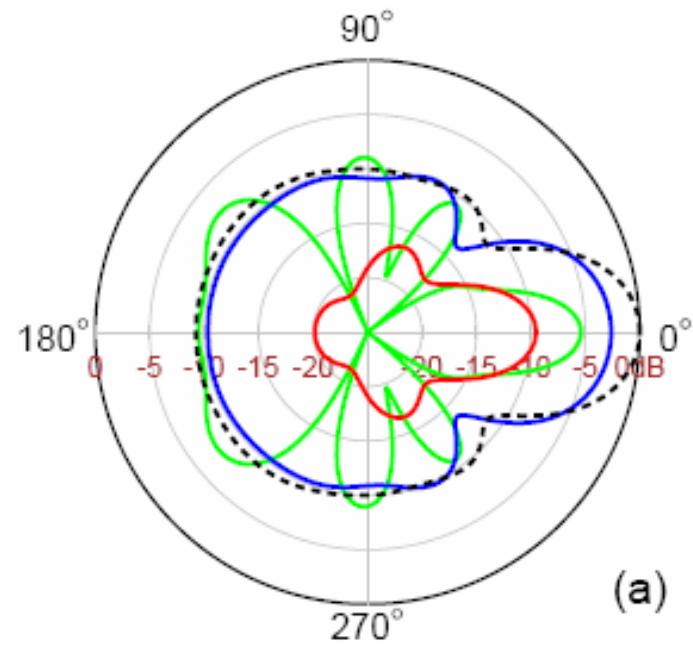
Improved simplification

Previous simplified



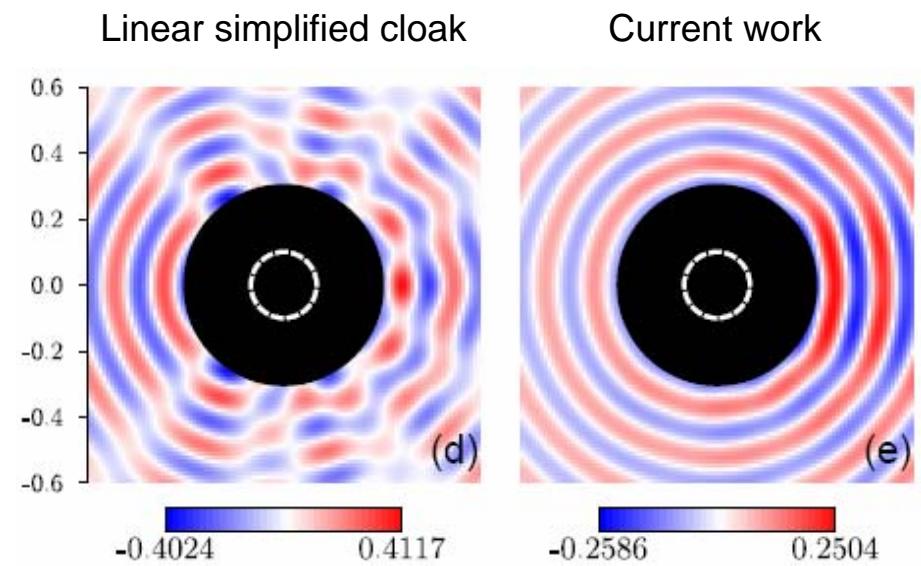


Improved simplification



(a)

- Bare cylinder
- Simplified linear cloak [2]
- Simplified quadratic cloak [3]
- Simplified linear cloak (improved)

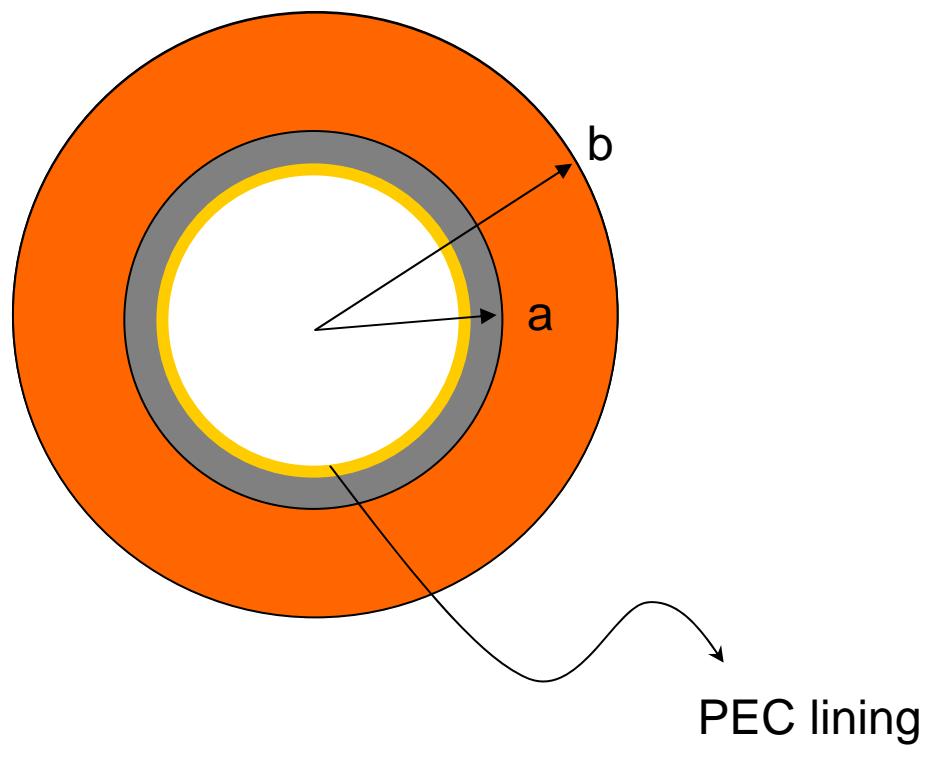


Scattered E_z field

1. M. Yan, Z. Ruan, and M. Qiu, Opt. Express **15**, p. 17772 (2007).
2. D. Schurig, et al., Science **314**, 977 (2006).
3. W. Cai et al., Appl. Phys. Lett. **91**, 111,105 (2007).



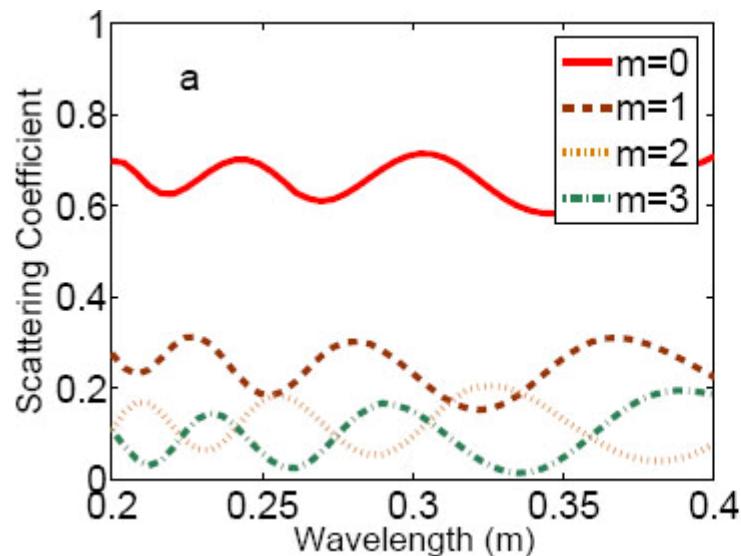
Can we cancel the zeroth order scattering?



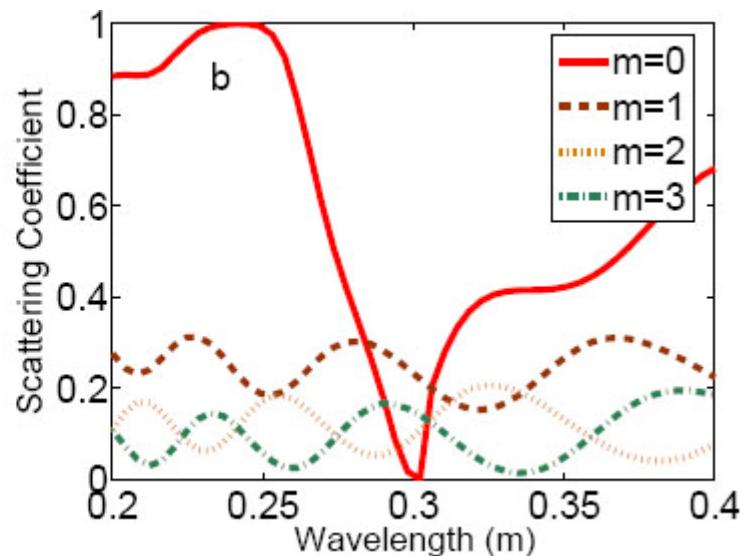


Cancellation of zeroth-order scattering

without monopole cancellation



with monopole cancellation

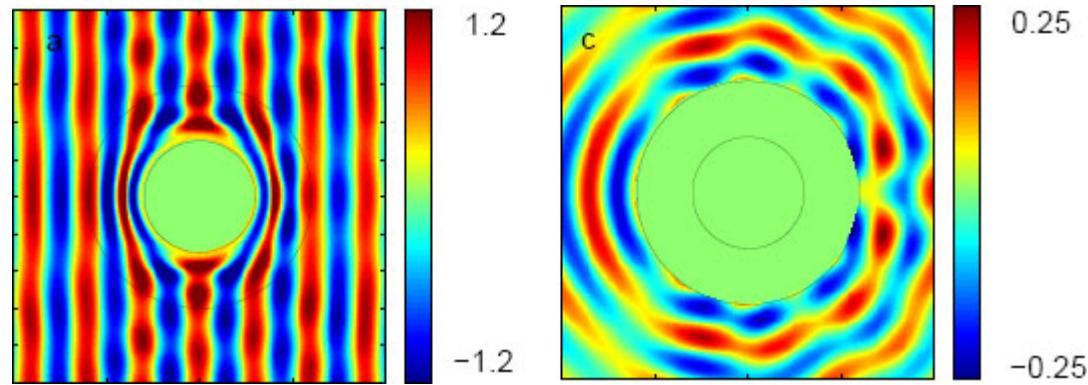


$$a = 0.3\text{m}, b = 0.6\text{m}$$

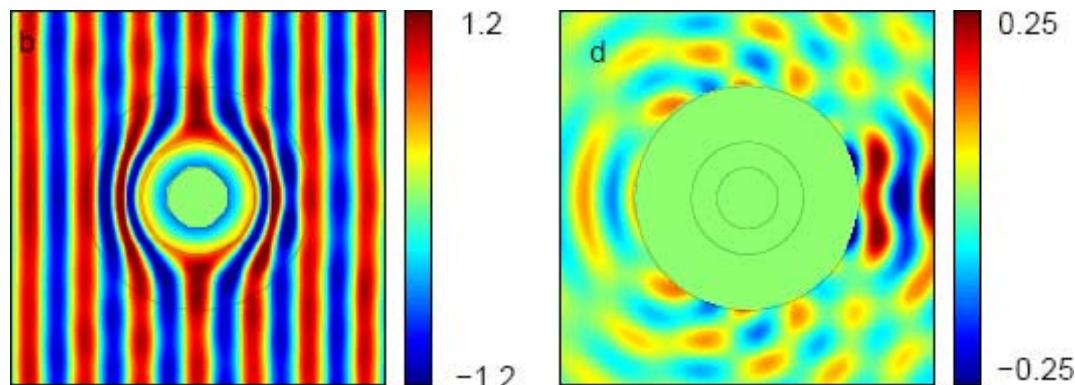


Cancellation of zeroth-order scattering

Without cancellation



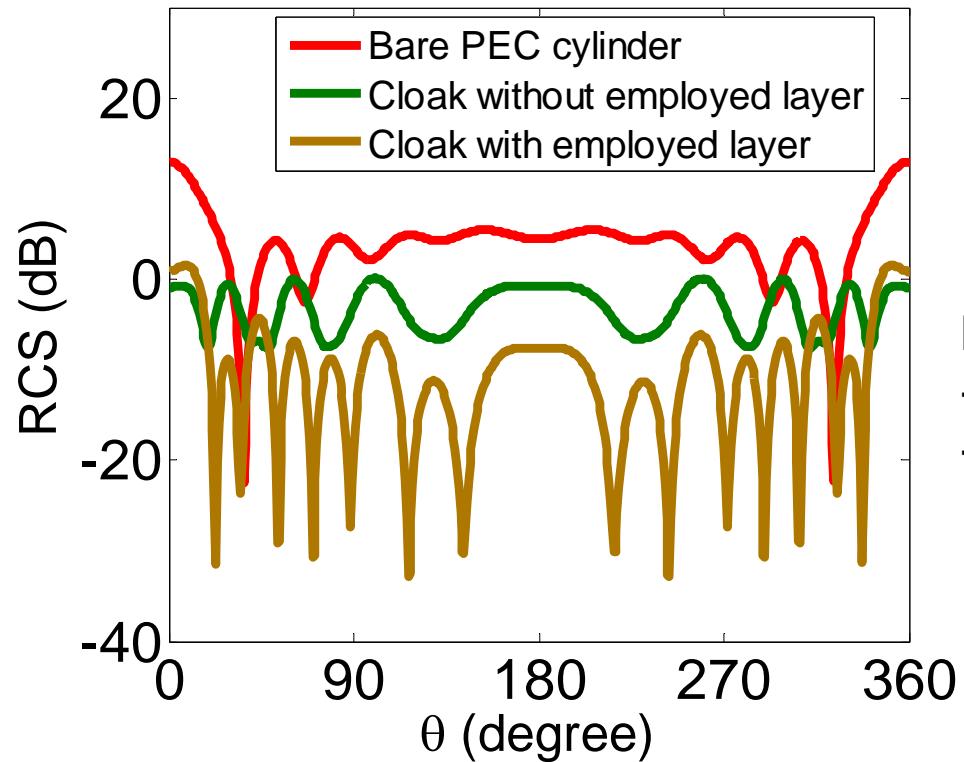
With cancellation



$$\lambda=0.3\text{m}$$



Far field radiation



Remarks:

- Mostly high order scatterings remain
- Wavelength dependent

Ref:

W. Yan, M. Yan, and Min Qiu, "Non-magnetic simplified cylindrical cloak with near perfect invisibility by suppressing zeroth order scattering," *Appl. Phys. Lett.*, in review.



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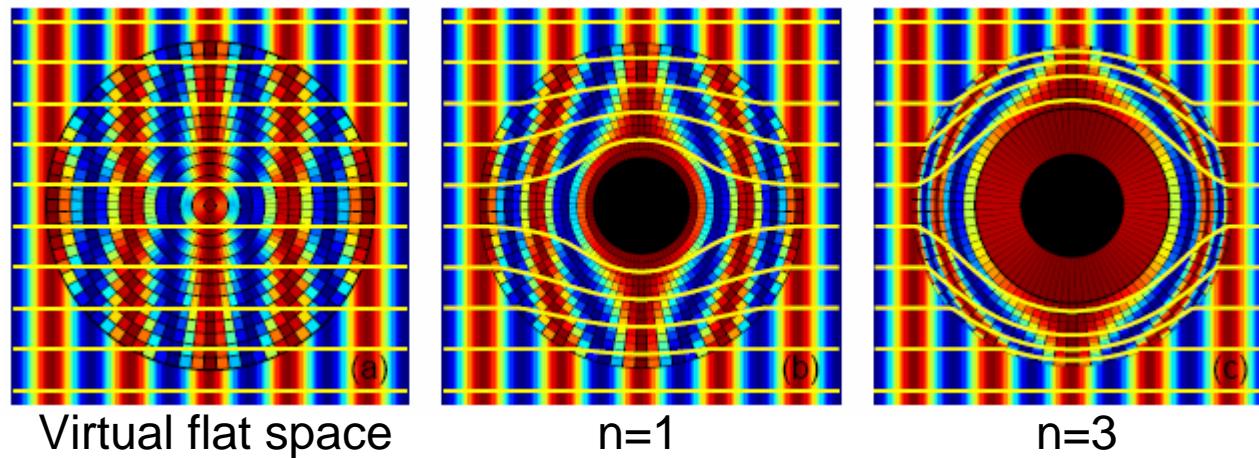


High-order cylindrical invisibility cloak

One possible class of transformation may take the form of

$$r' = \frac{b}{(b-a)^n} (r-a)^n,$$

n : transformation order.

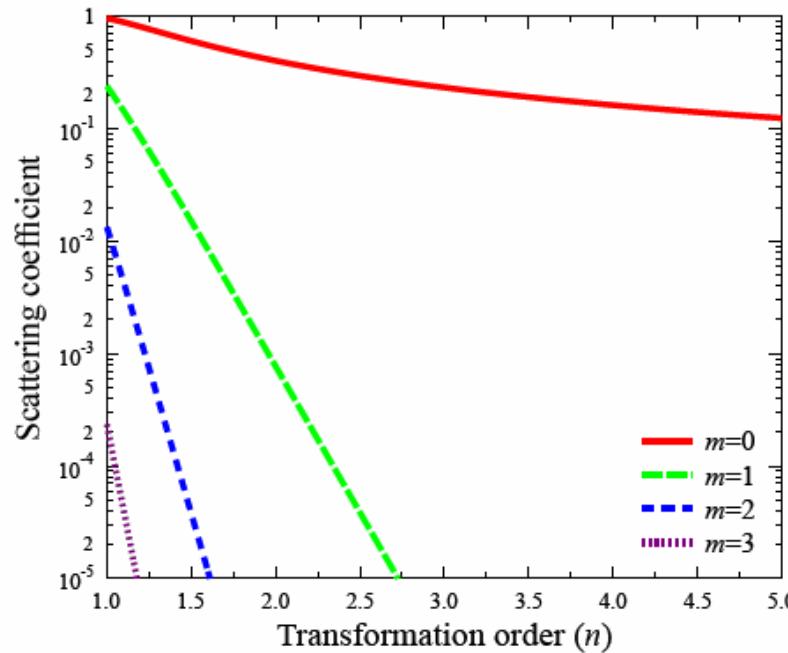


$$\varepsilon_r = \mu_r = \frac{r-a}{nr}, \quad \varepsilon_\theta = \mu_\theta = \frac{nr}{r-a}, \quad \varepsilon_z = \mu_z = \frac{nb^2(r-a)^{2n-1}}{(b-a)^{2n}r}.$$

Infinite parameters at $r=a$!



Effect of transformation order on scattering coefficients



$$a = 0.1\text{m}$$

$$b = 0.3\text{m}$$

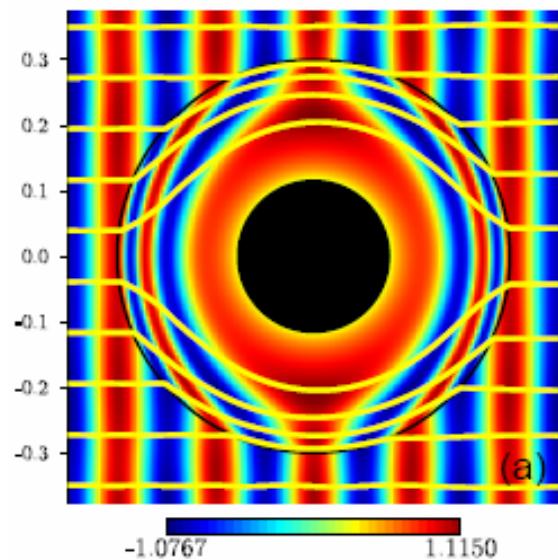
$$\lambda = 0.15\text{m}$$

Thickness of the layer peeled away: **$d=0.01\text{m}$**

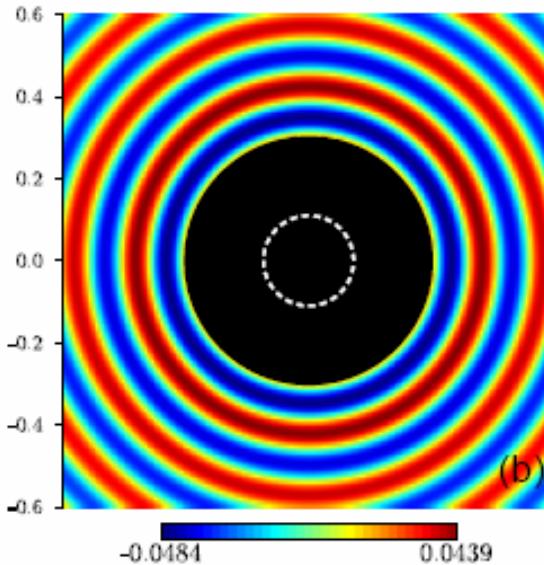
PEC lining is present



Field distributions



Ez



Scattered Ez

Transformation order used: n=3

Ref

M. Yan, W. Yan, L. Zhang, M. Qiu, Cylindrical Invisibility Cloak: Properties and Strategies for Practical Realization, NATO Advanced Research Workshop, Metamaterials for Secure Information and Communication Technologies, 7-10 May, 2008, Marrakech – Morocco.

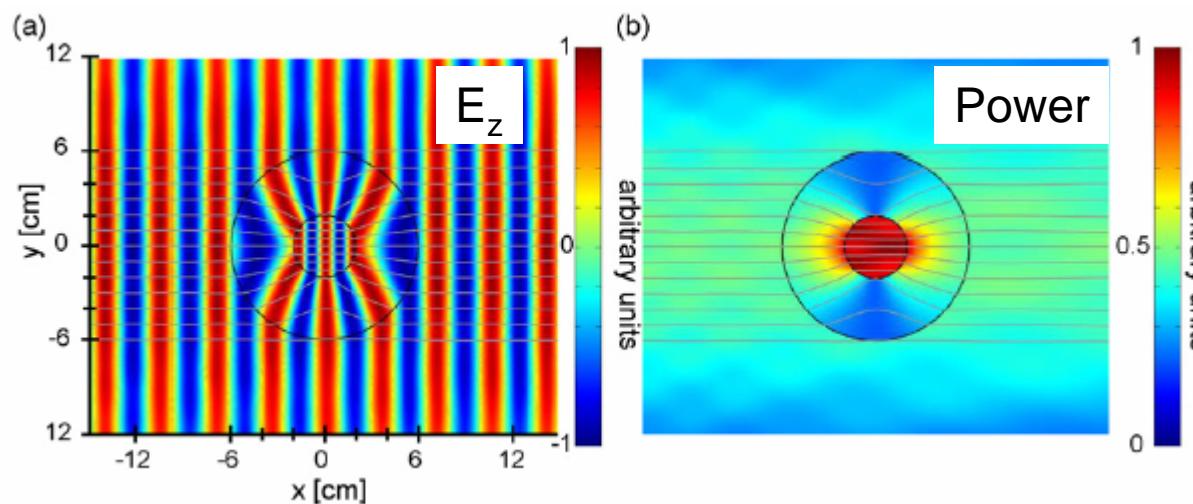
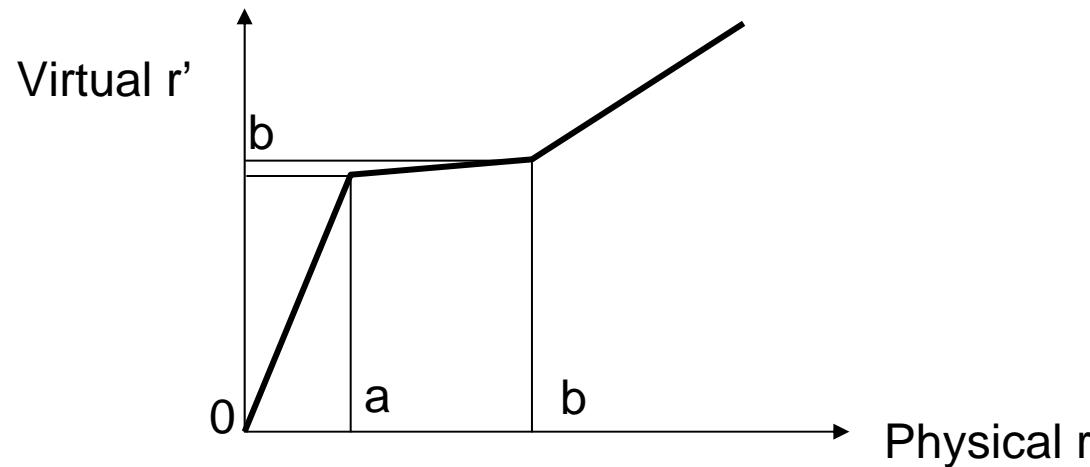


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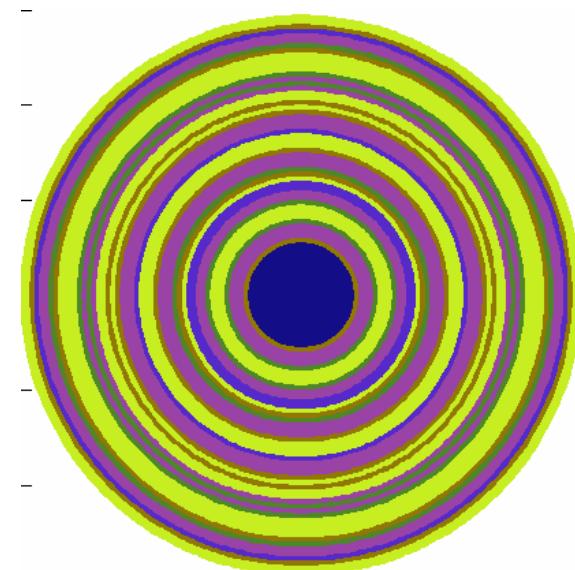
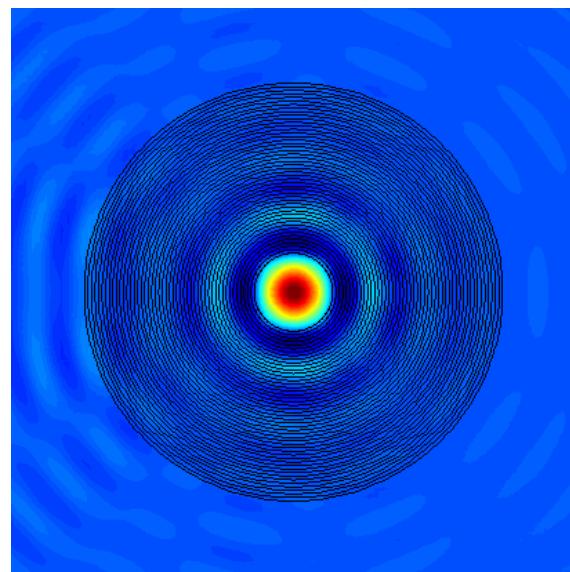
EM concentrator



Application:
● Photovoltaic fiber
● Fluid heating



Resonance based EM concentrator



- Enhancement factor 230
- Wavelength dependant



Resonance based EM concentrator

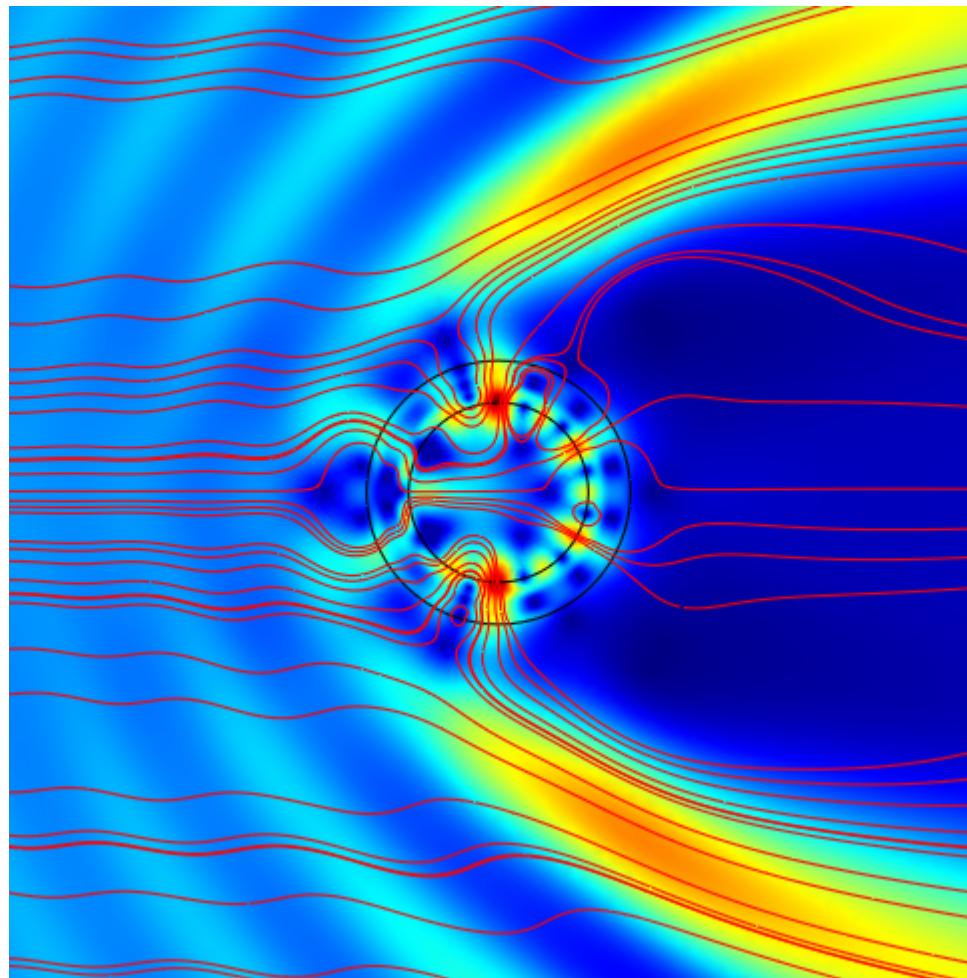


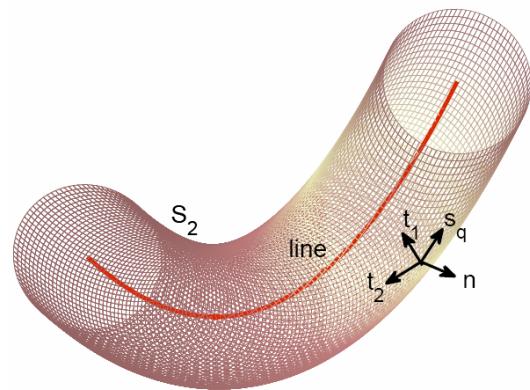


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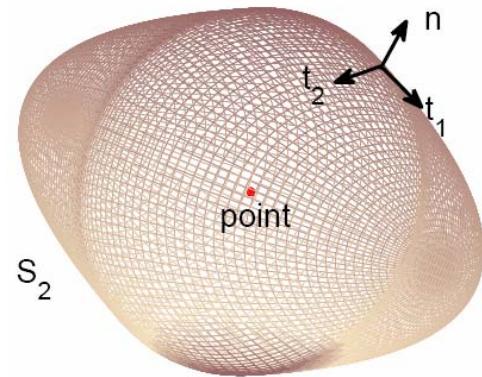
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Arbitrarily-shaped cloaks

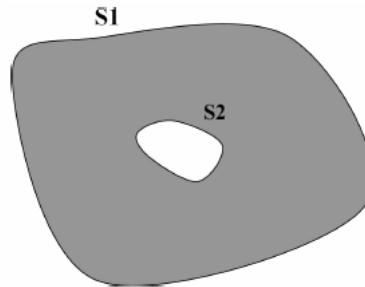


Blowing-up a line



Blowing-up a point

General cross-section:



Ref:

Wei Yan, Min Yan, Zhichao Ruan, Min Qiu, "Coordinate transformation makes perfect invisibility cloak with arbitrary shape," *New Journal of Physics*, vol. 10, 043040 (2008).



Arbitrarily-shaped cloak

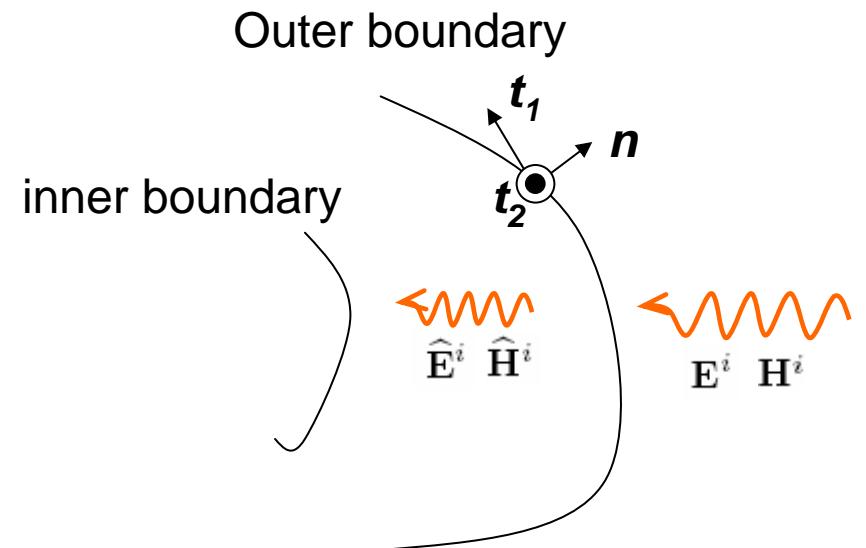
For both line-transformed or point-transformed cloaks,

$$\hat{\mathbf{E}}^i = \Lambda^T \mathbf{E}^i, \quad \hat{\mathbf{H}}^i = \Lambda^T \mathbf{H}^i,$$

At outer boundary:

$$\hat{E}_{t_1}^i = \mathbf{E}^i \cdot \hat{t}_1, \quad \hat{H}_{t_1}^i = \mathbf{H}^i \cdot \hat{t}_1,$$

$$\hat{E}_{t_2}^i = \mathbf{E}^i \cdot \hat{t}_2, \quad \hat{H}_{t_2}^i = \mathbf{H}^i \cdot \hat{t}_2,$$





Arbitrarily-shaped cloak

For line-transformed cloak,

$$\hat{E}_{t1}^i = \hat{H}_{t1}^i = 0.$$

However, the other components of fields are not zero. In particular,

$$\hat{E}_{t2}^i = (\hat{s} \cdot \hat{t}_2)[B_1, B_2, B_3]\mathbf{E}^i,$$

$$\hat{H}_{t2}^i = (\hat{s} \cdot \hat{t}_2)[B_1, B_2, B_3]\mathbf{H}^i,$$

$$\hat{E}_n^i = [F_1 + B_1(\hat{s} \cdot \hat{n}), F_2 + B_2(\hat{s} \cdot \hat{n}), F_3 + B_3(\hat{s} \cdot \hat{n})]\mathbf{E}^i,$$

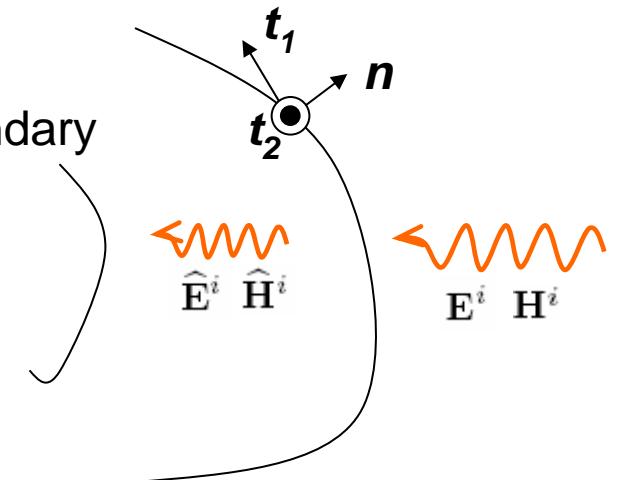
$$\hat{H}_n^i = [F_1 + B_1(\hat{s} \cdot \hat{n}), F_2 + B_2(\hat{s} \cdot \hat{n}), F_3 + B_3(\hat{s} \cdot \hat{n})]\mathbf{H}^i,$$

with

$$B_i = \sqrt{\partial b_i / \partial q_1^2 + \partial b_i / \partial q_2^2 + \partial b_i / \partial q_3^2}.$$

Outer boundary

inner boundary



Remark: Surface current will be induced at the inner surface.



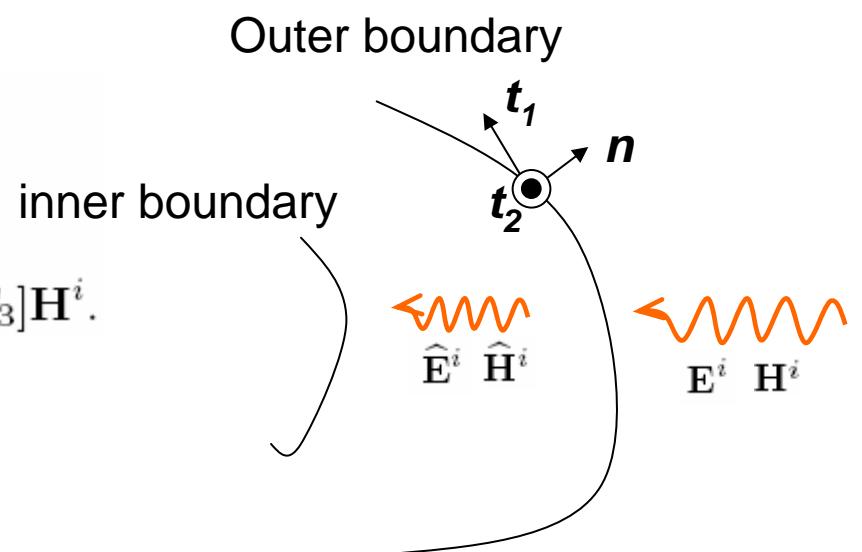
Arbitrarily-shaped cloak

For point-transformed cloak,

$$\hat{E}_{t_1}^i = \hat{H}_{t_1}^i = 0,$$

$$\hat{E}_{t_2}^i = \hat{H}_{t_2}^i = 0,$$

$$\hat{E}_n^i = [F_1, F_2, F_3] \mathbf{E}^i, \quad \hat{H}_n^i = [F_1, F_2, F_3] \mathbf{H}^i.$$



All tangential fields are zero.

Therefore no field discontinuity exists for perfect cloaking.



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Outcome

Papers:

1. Wei Yan, Min Yan, Zhichao Ruan, Min Qiu, "Coordinate transformation makes perfect invisibility cloak with arbitrary shape," *New Journal of Physics*, vol. 10, 043040 (2008).
2. M. Yan, W. Yan, L. Zhang, M. Qiu, *Cylindrical Invisibility Cloak: Properties and Strategies for Practical Realization*, NATO Advanced Research Workshop, Metamaterials for Secure Information and Communication Technologies, 7-10 May, 2008, Marrakech, Morocco.
3. W. Yan, M. Yan, and Min Qiu, "Non-magnetic simplified cylindrical cloak with near perfect invisibility by suppressing zeroth order scattering," *Appl. Phys. Lett.*, in review.

Code:

Matlab code for scattering calculation of multilayered cylindrical structures.



Thank you!

Questions?

